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N73-18928  
NASA CR-2174

NASA CR-2174

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FREE VIBRATIONS  
OF THERMALLY STRESSED  
ORTHOTROPIC PLATES WITH  
VARIOUS BOUNDARY CONDITIONS

*by Cecil D. Bailey and James C. Greetham*

*Prepared by*

THE OHIO STATE UNIVERSITY RESEARCH FOUNDATION

Columbus, Ohio

*for Langley Research Center*

NATIONAL AERONAUTICS AND SPACE ADMINISTRATION • WASHINGTON, D. C. • FEBRUARY 1973

1. Report No. <b>NASA CR-2174</b>	2. Government Accession No.	3. Recipient's Catalog No.	
4. Title and Subtitle <b>FREE VIBRATIONS OF THERMALLY STRESSED ORTHOTROPIC PLATES WITH VARIOUS BOUNDARY CONDITIONS</b>		5. Report Date <b>February 1973</b>	
		6. Performing Organization Code	
7. Author(s) <b>Cecil D. Bailey and James C. Greetham</b>		8. Performing Organization Report No.	
		10. Work Unit No.	
9. Performing Organization Name and Address  <b>The Ohio State University Research Foundation Columbus, Ohio</b>		11. Contract or Grant No. <b>NGR 36-008-109</b>	
		13. Type of Report and Period Covered	
12. Sponsoring Agency Name and Address <b>National Aeronautics and Space Administration Washington, D.C. 20546</b>		14. Sponsoring Agency Code	
15. Supplementary Notes			
16. Abstract <p>An analytical investigation of the vibrations of thermally stressed orthotropic plates in the prebuckled region is presented. The investigation covers the broad class of trapezoidal plates with two opposite sides parallel. Each edge of the plate may be subjected to different uniform boundary conditions. Variable thickness and arbitrary temperature distributions (analytical or experimental) for any desired combination of boundary conditions may be prescribed. Results obtained using this analysis are compared to experimental results obtained for isotropic plates with thermal stress, and to results contained in the literature for orthotropic plates without thermal stress. Good agreement exists for both sets of comparisons.</p>			
17. Key Words (Suggested by Author(s))  <b>Plate vibrations</b>  <b>Thermally stressed orthotropic plates</b>		18. Distribution Statement  <b>Unclassified - Unlimited</b>	
19. Security Classif. (of this report) <b>Unclassified</b>	20. Security Classif. (of this page) <b>Unclassified</b>	21. No. of Pages <b>100</b>	22. Price* <b>\$8.00</b>

## SUMMARY

An analytical investigation of the vibrations of thermally stressed orthotropic plates in the prebuckled region is presented. The investigation covers the broad class of trapezoidal plates with two opposite sides parallel. Each edge of the plate may be subjected to different uniform boundary conditions. Variable thickness and arbitrary temperature distributions (analytical or experimental) may be prescribed. Generality is achieved in the analysis through the treatment of boundary conditions, the choice of functions for stress distributions and deflection distributions, and the use of numerical integration for the evaluation of matrix elements. Results obtained using this analysis are compared to experimental results obtained for isotropic plates with thermal stress, and to results contained in the literature for orthotropic plates without thermal stress. Good agreement exists for both sets of comparisons. Calculations for several orthotropic plates with thermal stresses indicates that the effect of orthotropy on the frequencies may be large and should not be ignored.





## TABLE OF CONTENTS

<u>Section</u>	<u>Page</u>
Summary . . . . .	iii
List of Symbols . . . . .	vii
Note on Subscript Convention . . . . .	x
I. Introduction . . . . .	1
II. Application of the Energy Equation . . . . .	2
A. Potential Energy . . . . .	2
B. Complementary Energy . . . . .	5
C. The Equations in Matrix Form . . . . .	7
D. Deflection Function and Stress Function . . . . .	8
E. Boundary Conditions . . . . .	9
1. Displacement Function . . . . .	9
2. Stress Function . . . . .	11
III. Programming . . . . .	12
A. The Equations and the Plate Geometry . . . . .	12
B. Logic Flow Diagram . . . . .	13
C. Programming Boundary Conditions . . . . .	13
D. Numerical Integration . . . . .	15
E. Temperature Distribution . . . . .	17
F. Thickness Distribution . . . . .	19
G. Miscellaneous Comments . . . . .	19
IV. Results . . . . .	19
A. Comparison of Orthotropic Results Without Thermal Loading . . . . .	20

B.	Comparison of Isotropic Results With Thermal Loading . . . . .	21
C.	Comparison of Orthotropic Results With Thermal Loading . . . . .	22
D.	Effect of Stress Distribution . . . . .	23
V.	Concluding Remarks . . . . .	24
	References . . . . .	25

#### Appendix

A	Matrix Elements and Other Parameters . . . . .	47
B	Logic Flow Diagram . . . . .	49
C	Program Listing . . . . .	55

# LIST OF SYMBOLS

$[A]$	- stress matrix from complementary energy
$A_{pq,rs}$	- elements of $[A]$
$AR$	- aspect ratio, length squared/area
$a$	- plate length
$[a]$	- matrix of material constants in equation for strains in terms of stresses
$a_{11}, a_{12}, a_{22}, b_{12}$	- elements of $[a]$ , Equation (3)
$[B]$	- generalized stiffness matrix from bending energy
$B_{ij,kl}$	- element of $[B]$
$b$	- plate dimension measured along left edge from x-axis to top corner
$\bar{b}_1$	- plate dimension, plate width at left edge minus $b$
$b_1$	- ratio of $\bar{b}_1/b$
$b_i$	- $b_i(x,y) = 0$ , equation of $i^{th}$ portion of plate boundary
$C_{pq}$	- coefficient of the $pq^{th}$ term of the assumed stress function solution
$[E]$	- matrix of material constants in equation for stresses in terms of strains, $[E] = [a]^{-1}$
$E_{11}, E_{12}, E_{22}, G_{12}$	- elements of $[E]$
$F$	- stress function solution to the inplane equilibrium equations
$f_i$	- frequency of $i^{th}$ mode, cycles per second

$f$	- $f(x,y)$ - function which forces the assumed stress function solution to satisfy the stress boundary conditions
$g$	- $g(x,y)$ - forces the assumed displacement solution to satisfy the displacement boundary conditions
$h_{ij}$	- coefficient of the $ij^{\text{th}}$ term in the assumed displacement function solution
$h$	- $h(x,y)$ - function to represent any variation in plate thickness
$h_r$	- plate thickness at some reference point
$[M]$	- mid-plane energy matrix, associated with the thermal stresses moving through small out-of-plane displacements
$M_{ij,k\ell}$	- elements of $[M]$
$N_1, N_2, N_{12}$	- $\int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_1 dz, \int_{-\frac{h}{2}}^{\frac{h}{2}} \sigma_2 dz, \int_{-\frac{h}{2}}^{\frac{h}{2}} \tau_{12} dz$ respectively
$n$	- coordinate normal to plate boundary (in-plane)
$[T]$	- generalized mass matrix from kinetic energy
$T_{ij,k\ell}$	- element of $[T]$
$T$	- $T(x,y)$ - difference between the temperature at a point $(x,y)$ on the plate and the original, uniform reference temperature (non-dimensional)
$T_{\text{ref}}$	- difference between the temperature at some reference point on the plate and the original, uniform reference temperature
$\Delta T_{\text{cr}}$	- The magnitude of $T_{\text{ref}}$ at which the free vibration frequency vanishes. By definition, the thermal buckling temperature.



$t$	- time
$u, v, w$	- displacements in the x, y, and z directions respectively
$x, y$	- independent in-plane variables
$\bar{\alpha}$	- angle between plate leading edge and x-axis, measured positive counter-clockwise
$\alpha$	- taper parameter, $\alpha = \frac{a}{b} \tan \bar{\alpha}$
$\alpha_1, \alpha_2$	- coefficient of thermal expansion in x and y directions respectively
$ij$	- $ij^{th}$ term of general assumed displacement function
$\bar{\beta}$	- angle between x-axis and the line dividing the plate for thickness distribution purposes
$\beta$	- non-dimensional form of $\bar{\beta}$
$\{\Gamma\}$	- thermal loading matrix
$\Gamma_{rs}$	- element of $\{\Gamma\}$
$\bar{\gamma}$	- angle between plate trailing edge and x-axis, measured positive counter-clockwise
$\gamma$	- taper parameter, $\gamma = \frac{a}{b} \tan \bar{\gamma}$
$\gamma_{12}$	- shear strain
$\gamma_{pq}$	- $pq^{th}$ term of general assumed stress function
$\Delta T$	- an increment of $T_{ref}$ , gives magnitude of the temperature distribution under consideration
$\epsilon_1, \epsilon_2$	- normal strains in the x and y directions respectively
$\eta$	- non-dimensional independent space variable, $\eta = y/b$

$\lambda_i$	- vibration eigenvalue, $\lambda_i = \omega_i \frac{a^2}{h_r} \sqrt{12\rho/E_{11}}$
$\xi$	- non-dimensional independent space variable, $\xi = x/a$
$\pi$	- energy of a system per unit time
$\pi^*$	- complimentary energy
$\rho$	- plate material density, mass per unit volume
$\sigma_1, \sigma_2$	- normal stresses in x and y directions respectively
$\tau_{12}$	- shear stress
$\omega$	- vibration frequency of the thermally stressed plate, radians/sec.
$\omega_0$	- vibration frequency of the plate at $T = 0$ , radians/sec.

#### NOTE ON SUBSCRIPT CONVENTION

Numeric subscripts indicate the component of a quantity in a coordinate direction (e.g.,  $\sigma_1$  - normal stress in the 1 or x - direction). A subscript of x, y,  $\xi$ , or  $\eta$  denotes differentiation with respect to that independent variable (e.g.,  $(\sigma_1)_x = \frac{\partial}{\partial x} (\sigma_1)$ ). All other alphabetic subscripts (i, j, k, , p, q, etc.) will refer to either terms in a series or elements in a matrix.

## I. INTRODUCTION

Considerable work has been reported in the literature on the problem of finding the frequencies and modes of vibration of a rectangular orthotropic plate at ambient temperature. A combination of the work of Hearmon (Ref. 1, 2, 3); Hoppman, Huffington, and Magness (in various combinations, Ref. 4, 5, 6, 7, 8); Kanazawa and Kawai (Ref. 9) and Wah (Ref. 10) provide solutions for rectangular plates with any boundary condition except completely free.

In contrast, no literature was found pertaining to the free vibration frequencies of an orthotropic plate subjected to a thermal loading. For the special case of a thermally stressed isotropic plate, the torsion mode of the plate with cantilever boundary conditions has been rather thoroughly investigated (Ref. 11, 12, 13, 14, 15).

Ref. 16 presents an analysis of thermally stressed isotropic plates for various boundary conditions, ranging from plates completely clamped through several combinations of mixed boundary conditions to plates with all edges completely free. This paper extends the analysis of Ref. 16 to include orthotropic plates with a thorough discussion of the associated computer program.

In the sense that both compatibility and equilibrium are satisfied as closely as one pleases at every point interior to the plate and on the boundary, this paper presents an analysis that provides, in a practical computational sense, a solution to the thermally stressed plate vibration problem for all trapezoidal plates with two opposite sides parallel, and with one of the axes of elastic symmetry parallel to these sides, restrictions that could be easily relaxed.

The analysis and associated computer program are of sufficient generality that isotropic plates are included as a special case of orthotropic plates. Various boundary conditions may be arbitrarily assigned to the different sides of the plate. Thus, the solution for the vibrations of thermally stressed plates with boundary conditions ranging from completely clamped to completely free with any combination of clamped, pinned, and/or free edges may be obtained.

A small number of quantitative strain measurements (not included herein) plus the abundance of experimental dynamic response data for various planform shapes and boundary conditions of isotropic plates indicates that the stress distributions

as determined herein are correct.

No thermally stressed orthotropic plate data, either analytical or experimental, were found in the literature; however, for orthotropic plates without thermal stress, comparison is made to both analytical and experimental data from the literature. Further comparison is made with experimental data for several modes of thermally stressed isotropic plates with various planform shapes, boundary conditions, and temperature distributions.

## II. APPLICATION OF THE ENERGY EQUATION

Because of a difference in notation between that used herein and that used in other sources, the derivation of the expressions for potential and complementary energy will be shown. Except for this difference, the procedures used are well known. Additional details, if desired, may be found in Ref. 17.

### A. Potential Energy

The forces are taken as the independent variables and the variation of the total energy is taken with respect to the displacements. With the assumptions of plane stress and no body or surface forces,

$$\delta\pi = \iiint \{(\sigma_1 \delta\epsilon_1 + \sigma_2 \delta\epsilon_2 + \tau_{12} \delta\gamma_{12}) + \rho \ddot{w} \delta w\} dx dy dz \quad (1)$$

The orthotropic stress-strain relations will be taken as,

$$\{\epsilon\} = [a] \{\sigma\} + \{\alpha\} T \quad (2)$$

where



$$[a] = \begin{bmatrix} a_{11} & a_{12} & 0 \\ a_{12} & a_{22} & 0 \\ 0 & 0 & b_{12} \end{bmatrix} ; \quad \{\alpha\} = \begin{Bmatrix} \alpha_1 \\ \alpha_2 \\ 0 \end{Bmatrix} \quad (3)$$

$$\{\epsilon\} = \begin{Bmatrix} \epsilon_1 \\ \epsilon_2 \\ \gamma_{12} \end{Bmatrix} \quad \text{and} \quad \{\sigma\} = \begin{Bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{Bmatrix}$$

The inverse of eq. (2) is

$$\{\sigma\} = [E] \{\epsilon\} - [E] \{\epsilon\}^T \quad (4)$$

The Von Karman strain-displacement equations are used,

$$\begin{aligned} \epsilon_1 &= u_x + \frac{1}{2} (w_x)^2 - z w_{xx} \\ \epsilon_2 &= v_y + \frac{1}{2} (w_y)^2 - z w_{yy} \\ \gamma_{12} &= u_y + v_x + w_x w_y - 2z w_{xy} \end{aligned} \quad (5)$$

Substitute eqs. 4 and 5 into eq. 1, carry out the indicated operations, neglect fourth order terms, and integrate through the thickness to get,

$$\begin{aligned}
\delta \pi = & \frac{1}{2} \iint \{ N_1 u_x + N_2 v_y + N_{12} (u_y + v_x) \\
& + \frac{h^3}{12} [E_{11} (w_{xx})^2 + E_{22} (w_{yy})^2 + 2E_{12} w_{xx} w_{yy} + 4G_{12} (w_{xy})^2] \\
& + N_1 (w_x)^2 + N_2 (w_y)^2 + 2N_{12} w_x w_y \\
& + 2 \rho h \ddot{w} \} dx dy = 0
\end{aligned} \tag{6}$$

where,

$$\begin{aligned}
N_1 &= \int_{-\frac{h}{2}}^{\frac{h}{2}} [E_{11} u_x + E_{12} v_y - T (E_{11} \alpha_1 + E_{12} \alpha_2)] dz \\
N_2 &= \int_{-\frac{h}{2}}^{\frac{h}{2}} [E_{12} u_x + E_{22} v_y - T (E_{12} \alpha_1 + E_{22} \alpha_2)] dz
\end{aligned} \tag{7}$$

$$N_{12} = \int_{-\frac{h}{2}}^{\frac{h}{2}} G_{12} (u_y + v_x) dz,$$

Taking the variation with respect to  $u$  gives,

$$\iint (N_1 \delta u_x + N_{12} \delta u_y) dx dy = 0.$$

Integrating this result by parts, noting that for any solution to a particular problem the boundary conditions must be satisfied, leaves the in-plane equilibrium equation in the  $x$ -direction,

$$(N_1)_x + (N_{12})_y = 0 \tag{8a}$$

Performing a similar series of operations on  $v$  gives for the  $y$ -direction,

$$(N_{12})_x + (N_2)_y = 0 \quad (8b)$$

These eqs. (8) have the solution,  $F$ , such that,

$$\begin{aligned} N_1 &= F_{yy} \\ N_2 &= F_{xx} \\ N_{12} &= -F_{xy} \end{aligned} \quad (9)$$

Thus the variational expression for potential energy of a thermally stressed, orthotropic plate becomes,

$$\begin{aligned} \delta\pi = \frac{1}{2} \iiint \left\{ \frac{h^3}{12} [E_{11}(w_{xx})^2 + E_{22}(w_{yy})^2 + 2E_{12}w_{xx}w_{yy} \right. \\ \left. + 4G_{12}(w_{xy})^2] \right. \\ \left. + [F_{yy}(w_x)^2 + F_{xx}(w_y)^2 - 2F_{xy}w_xw_y] \right. \\ \left. + 2ph \ddot{w} w \right\} dx dy = 0 \end{aligned} \quad (10)$$

#### B. Complementary Energy

In developing the expression for complementary energy, the unknown forces are varied and the displacements are held constant. Thus, with the same assumptions as used in the treatment of potential energy,

$$\delta\pi^* = \iiint \{ \epsilon_1 \delta\sigma_1 + \epsilon_2 \delta\sigma_2 + \gamma_{12} \delta\tau_{12} \} dx dy dz - \delta w_B = 0$$

where  $W_B$  represents the work done by the stresses on the portion of the boundary on which the displacements are specified. In this treatment, if the displacements on any part of the boundary of the plate are to be specified, they will be specified to be zero. Thus  $W_B = 0$ .

Substitute eq. (4) for the stresses to get,

$$\begin{aligned}\delta\pi^* = & \iiint \{ \epsilon_1 \delta [E_{11} \epsilon_1 + E_{12} \epsilon_2 - (E_{11} \alpha_1 + E_{12} \alpha_2) T] \\ & + \epsilon_2 \delta [E_{22} \epsilon_2 + E_{12} \epsilon_1 - (E_{12} \alpha_1 + E_{22} \alpha_2) T] \\ & + \gamma_{12} \delta G_{12} \gamma_{12} \} dx dy dz = 0\end{aligned}$$

Because the bending stresses have been expressed in terms of the displacement only and small deflections are assumed, the stresses that remain in the equations are not functions of the out-of-plane displacements; they are "membrane stresses" resulting only from the in-plane displacements and/or temperature. Thus, only the linear strains resulting from in-plane deformation need to be considered and eq. (5) can be simplified to,

$$\begin{aligned}\epsilon_1 &= u_x \\ \epsilon_2 &= v_y \\ \gamma_{12} &= (u_y + v_x)\end{aligned}$$

With these relations, the complementary energy can be expressed as,

$$\begin{aligned}\delta\pi^* = & \iiint \{ u_x \delta [E_{11} u_x + E_{12} v_y - (E_{11} \alpha_1 + E_{12} \alpha_2) T] \\ & + v_y \delta [E_{12} u_x + E_{22} v_y - (E_{12} \alpha_1 + E_{22} \alpha_2) T] \\ & + (u_y + v_x) \delta [G_{12} (u_y + v_x)] \} dx dy dz = 0\end{aligned}$$



Integrate through the thickness, substitute equations (7) and (9) and define the strains to be,

$$\epsilon_1 = u_x = \frac{1}{h} (a_{11} F_{yy} + a_{12} F_{xx}) + \alpha_1 T$$

$$\epsilon_2 = v_y = \frac{1}{h} (a_{22} F_{xx} + a_{12} F_{yy}) + \alpha_2 T$$

$$\gamma_{12} = u_y + v_x = -\frac{1}{h} b_{12} F_{xy} ,$$

from which the complementary energy for a thermally stressed, orthotropic plate becomes,

$$\begin{aligned} \delta\pi^* = \delta \iint \{ & \frac{1}{2h} [a_{11} (F_{yy})^2 + a_{22} (F_{xx})^2 + 2a_{12} F_{xx} F_{yy} \\ & + b_{12} (F_{xy})^2] \end{aligned} \quad (11)$$

$$+ (\alpha_1 F_{yy} + \alpha_2 F_{xx}) T \} dx dy = 0 .$$

### C. The Equations in Matrix Form

Consider first the potential energy, eq. (10). Assume a displacement function of the form,

$$w(x,y,t) = \sum_{i=0}^N \sum_{j=0}^M h_{ij}(t) \alpha_{ij}(x,y) \quad (12)$$

where each  $\alpha_{ij}(x,y)$ , (1) satisfies the displacement boundary conditions, (2) is continuous, and (3) has at least continuous first derivatives.

Substitute this into eq. (10), take the variation with respect to  $h_{kl}$  and collect coefficients of like  $h_{ij}$  to get the matrix equation,

$$[B] \{h_{ij}\} + K_1 [M] \{h_{ij}\} - \lambda^2 [T] \{h_{ij}\} = 0 \quad (13)$$

where the non-dimensionalized matrix elements and associated parameters are given in Appendix A.

Now assume a stress function of the form,

$$F(x,y) = \sum_{p=0}^S \sum_{q=0}^T C_{pq} \gamma_{pq}(x,y) \quad (14)$$

where each  $\gamma_{pq}(x,y)$ , (1) satisfies the stress boundary conditions, (2) is continuous, and (3) has at least continuous first derivatives.

Substitute this into eq. (11), take the variation with respect to  $C_{rs}$  and collect coefficients of like  $C_{pq}$  to get the matrix equation,

$$[A] \{\hat{C}_{pq}\} + K_2 \{\Gamma\} = 0 \quad (15)$$

where the matrix elements are also given in Appendix A.

Thus, given a temperature distribution,  $\{\Gamma\}$  can be calculated, eq. (15) can be solved for  $\{\hat{C}_{pq}\}$ , and values of the derivatives of the stress function can be found. Using this information, the elements,  $M_{ij,kl}$ , can be calculated and eq. (13) can be solved for the vibration frequencies and modes with the buckling mode and  $\Delta T_{cr1}$  obtainable as a limiting case when  $\lambda_1^2 = 0$ .

#### D. Deflection Function and Stress Function

At this point, a choice will be made concerning the form of the assumed deflection function and stress function. By observing the physical system, it can be seen that the deflected surface of the plate and the stresses within the plate are continuous and have at least continuous first derivatives. Thus, the functions to be assumed as solutions to the problem must belong to the class of functions which are continuous and have at least continuous first derivatives. The assumed solution must also satisfy the boundary conditions discussed in the next section.

A truncated power series in the independent space variables satisfies the continuity requirements. Thus, the functions assumed for the deflection,  $w$ , and for the stress function,  $F$ , will be truncated power series.

#### E. Boundary Conditions

The polynomial resulting from a truncated power series will not in general satisfy the boundary conditions. Therefore, the polynomial representation must be modified by an additional function which forces satisfaction of the required boundary conditions.

Let the displacement function have the form,

$$w(x,y) = g(x,y) + \sum_{i=0}^N \sum_{j=0}^M h_{ij} x^i y^j$$

where  $g(x,y)$  is the boundary condition function which insures satisfaction of the displacement conditions at the boundary. The stress function will have the form,

$$F(x,y) = f(x,y) + \sum_{p=0}^S \sum_{q=0}^T C_{pq} x^p y^q$$

where  $f(x,y)$  is the boundary condition function which insures satisfaction of equilibrium in the plane of the plate at the boundary, i.e., the stress boundary condition. The specific form of each will now be considered.

##### 1. Displacement Function

Three types of displacement boundary conditions are considered herein:

- (a) Both displacement and slope normal to the edge of the plate are assumed to be zero; i.e., the edge is clamped.
- (b) Only the displacement is assumed to be zero and the slope is left unspecified resulting in a pinned (simply-supported) condition.
- (c) Both slope and displacement are left unspecified leaving the edge completely free.

Now, given a particular plate geometry, the equation of the boundary may be expressed as a polynomial, say,

$$b(x,y) = 0 \quad .$$

Therefore, in order to force the displacement to be zero on the boundary, simply let,

$$g(x,y) = b(x,y) \quad ,$$

so that for any point on the boundary,  $(x_B, y_B)$ , the deflection will be

$$w(x_B, y_B) = g(x_B, y_B) \sum_{i=0}^N \sum_{j=0}^M h_{ij} x_B^i y_B^j = 0$$

This satisfies condition (b) because the first derivative,  $\frac{\partial w}{\partial n}$ , will not in general be zero but will be left to take on whatever value is required for a minimum energy configuration.

Condition (a) may be satisfied by letting

$$g(x,y) = [b(x,y)]^2$$

The displacement will again be zero, but now the first derivative will also be zero on the boundary:

$$\begin{aligned} \frac{\partial w}{\partial n} &= \sum_{i=0}^N \sum_{j=0}^M h_{ij} (g(x,y) \frac{\partial (x^i y^j)}{\partial n} + x^i y^j \frac{\partial g(x,y)}{\partial n}) \\ &= \sum_{i=0}^N \sum_{j=0}^M h_{ij} \{ [b(x,y)]^2 \frac{\partial}{\partial n} (x^i y^j) + 2 x^i y^j b(x,y) \frac{\partial}{\partial n} b(x,y) \} \end{aligned}$$

and at a point  $(x_B, y_B)$  on the plate boundary,

$$\frac{\partial w}{\partial n} = 0 \quad .$$



Case (c) may be satisfied simply by letting

$$g(x,y) \equiv 1 = [b(x,y)]^0$$

Thus, in the case of an edge free to displace out of the plane of the plate, both the displacement and slope will be left to take on whatever values are required for a minimum energy configuration.

If the plate is a polygon of N sides, write

$$g(x,y) = \prod_{i=1}^N [b_i(x,y)]^{k_i},$$

where  $b_i(x,y)$  is the equation of the  $i^{\text{th}}$  side of the polygon. (The sides need not be straight.)  $k_i$  will be either 0, 1, or 2 as described above.

## 2. Stress Function

Two types of stress boundary conditions are considered herein:

- (a) The in-plane stresses normal to a boundary are specified to be zero. That is, the plate is left free to expand in the in-plane direction.
- (b) The stresses are completely unspecified or, equivalently, the in-plane displacements are specified to be zero. The stresses will take on whatever values are required for satisfaction of equilibrium.

Thus, condition (a) will be termed "free" and (b) will be termed "clamped". These conditions are fulfilled in a fashion similar to that used with the deflection function.

Recall from classical elasticity theory that the stresses normal to the edge of a plate will be zero if the stress function and its first derivative (normal to the edge) vanish there. Therefore, on any portion of the boundary on which a free condition is desired, the equilibrating function is,

$$f(x,y) - [b_i(x,y)] = 0$$

where, as before,  $b_i(x,y) = 0$  is the equation of that portion of the boundary.

A clamped condition on any portion of the boundary can be satisfied by setting,

$$f(x,y) = [b_i(x,y)]^0 \equiv 1 .$$

Thus, as was done with the boundary condition function, the equilibrating function is

$$f(x,y) = \prod_{i=1}^N [b_i(x,y)]^{k_i}$$

where  $k_i = 0, 2$  for clamped or free conditions respectively on the "ith" side of the polygon.

With these conditions, it is now possible to specify six different types of boundary conditions on any plate edge. The first letter of the notation used herein will denote the displacement boundary condition by using F = free, P = pinned or simply supported, and C = clamped. The second letter denotes the stress condition so that designations possible for any given edge are:

DESIGNATION	DISPLACEMENT	STRESS
(1) F - F	free	free
(2) P - F	pinned	free
(3) C - F	clamped	free
(4) F - C	free	clamped
(5) P - C	pinned	clamped
(6) C - C	clamped	clamped

### III. PROGRAMMING

#### A. The Equations and the Plate Geometry

The equations to be programed are eqs. 13 and 15 with the matrix elements as given in Appendix A.

The planform and descriptive parameters of the plates considered herein are shown in Fig. 1. The restriction that two of the sides are parallel was made to simplify the numerical integration scheme.

The plate edges are numbered clockwise (in the top view) beginning with the edge containing the origin. Thus, the equations of the four edges as used in the boundary condition functions are:

$$\begin{aligned}
 b_1(\xi, \eta) &= \xi = 0 \\
 b_2(\xi, \eta) &= (1 + \gamma\xi - \eta) = 0 \\
 b_3(\xi, \eta) &= (1 - \xi) = 0 \\
 b_4(\xi, \eta) &= (b_1 - \gamma\xi + \eta) = 0
 \end{aligned}
 \tag{16}$$

#### B. Logic Flow Diagram

The organization of the parts of the program is presented in a logic flow diagram shown in Appendix B. It should be noted first that if several sets of material properties and/or aspect ratios are to be investigated, the [B] and [A] matrices need not be integrated each time if the integrands are treated as four separate terms. Each of these terms need only be integrated once, then multiplied by the appropriate constant and added together to make up the whole integral for either isotropic or orthotropic materials. The [T] and [M] matrices are independent of both the material properties and aspect ratio. The program is structured to include this feature. A complete listing of the program is given in Appendix C.

#### C. Programming Boundary Conditions

The subroutine "FUNCTN" which calculates the displacement boundary condition function and the equilibrating function and their derivatives is very straightforward. Since both functions have the same form, the same subroutine can be used to simplify the user's task of calculating the exponents required. There is no need, for example, to remember that a zero exponent on the stress function means a clamped edge while the same exponent in the displacement boundary condition function means a free edge.

The first step was to write down the function and its five derivatives, leaving the exponents as variables. For example,

$$F = \xi^{I_1} (1+\alpha\xi-\eta)^{I_2} (1-\xi)^{I_3} (b_1-\gamma\xi+\eta)^{I_4},$$

$$\frac{\partial F}{\partial \eta} = -I_2 [\xi^{I_1} (1+\alpha\xi-\eta)^{I_1-1} (1-\xi)^{I_3} (b_1-\gamma\xi+\eta)^{I_4}]$$

$$+ I_4 [\xi^{I_1} (1+\alpha\xi-\eta)^{I_2} (1-\xi)^{I_3} (b_1-\gamma\xi+\eta)^{I_4-1}]$$

$$\partial^2 F / \partial \eta^2 = I_2 (I_2-1) [\xi^{I_1} (1+\alpha\xi-\eta)^{I_2-2} (1-\xi)^{I_3} (b_1-\gamma\xi+\eta)^{I_4}]$$

$$+ I_4 (I_4-1) [\xi^{I_1} (1+\alpha\xi-\eta)^{I_2} (1-\xi)^{I_3} (b_1-\gamma\xi+\eta)^{I_4-2}]$$

$$- 2I_2 I_4 [\xi^{I_1} (1+\alpha\xi-\eta)^{I_2-1} (1-\xi)^{I_3} (b_1-\gamma\xi+\eta)^{I_4-1}]$$

Thus, it can be seen that  $I_i$ ,  $I_i-1$ , and  $I_i-2$  are required for calculating the function and its derivatives. In the subroutine, the variable IEX(I,J) contains these quantities. The "I" refers to the four factors making up the function and "J" to  $I_i-0$ ,  $I_i-1$ , or  $I_i-2$ . This is done in "DO-LOOP" number five.

Next, this information is used to calculate all the factors T (M,K) required for the function and its derivatives. For example,

$$T(1,1) = \xi^{I_1}$$

$$T(1,2) = \xi^{I_1-1}$$

$$T(3,1) = (1-\xi)^{I_3}$$

$$T(4,3) = (b_1+\gamma\xi-\eta)^{I_4-2}$$

Finally, this information is used to calculate the function and its derivatives.

#### D. Numerical Integration

Integration of the elements of the various matrices is performed using the Gaussian Quadrature rule. The plate is divided into two parts by the line at angle  $\beta$ . This provides for more accurate results when the leading or forward part of the plate has a different thickness function than does the rearward part. Since the Gaussian Quadrature is defined on the interval  $[-1, 1]$  it is necessary to transform the points and coordinates.

$$\text{If } u = \phi(x) \quad \text{then } \int_a^b f(u) du = \int_{-1}^1 f[\phi(x)] \frac{d\phi(x)}{dx} dx$$

Then if

$$\int_{-1}^1 f(x) dx \approx \sum_{k=1}^N w_k f(x_k) ,$$

$$\int_a^b f(u) du \approx \sum_{k=1}^N W_k F(u_k)$$

where

$$W_k = \frac{d\phi(x)}{dx} w_k, \quad u_k = \phi(x_k).$$

For this problem,

$$\begin{aligned} & \int_0^1 \left\{ \int_a^b f(\xi, \eta) d\eta + \int_c^a f(\xi, \eta) d\eta \right\} d\xi \\ &= \int_{-1}^1 \left\{ \int_{-1}^1 f[\phi(x, y), \psi_1(x, y)] \frac{\partial \psi_1}{\partial y} dy \right. \\ & \quad \left. + \int_{-1}^1 f[\phi(x, y), \psi_2(x, y)] \frac{\partial \psi_2}{\partial y} dy \right\} \frac{\partial \phi}{\partial x} dx , \end{aligned}$$

where

$$a = \beta\xi, \quad b = 1+\alpha\xi, \quad c = \gamma\xi - b_1$$

Thus

$$\xi = 1/2 (x+1) \quad (17)$$

$$\eta = 1/2 \{y[1+ (\alpha-\beta)\xi] + 1 + (\alpha+\beta)\xi\}; \quad (\beta\xi \leq \eta \leq 1+\alpha\xi)$$

$$\eta = 1/2 \{y[b_1 + (\beta-\gamma)\xi] - b_1 + (\beta+\gamma)\xi\}; \quad (\gamma\xi \leq \eta \leq \beta\xi)$$

Hence,

$$\frac{\partial \phi}{\partial x} = 1/2$$

$$\frac{\partial \psi_1}{\partial y} = 1/2 [1+(\alpha-\beta)\xi]$$

$$\frac{\partial \psi_2}{\partial y} = 1/2 [b_1 + (\beta-\gamma)\xi]$$

and the integrals may be evaluated by

$$\begin{aligned} & \int_0^1 \left\{ \int_a^b f(\xi, \eta) d\eta + \int_c^a f(\xi, \eta) d\eta \right\} d\xi \\ & \approx \sum_{k=1}^N \frac{w_k}{2} \left\{ \sum_{L=1}^N \frac{w_L}{2} [b_1 + (\beta-\gamma)\xi_k] f(\xi_k, \eta_L) \right. \\ & \quad \left. + \sum_{M=N+1}^{2N} \frac{w_M}{2} [1 + (\alpha-\beta)\xi_k] f(\xi_k, \eta_M) \right\} \end{aligned}$$

where  $w_k$ ,  $w_L$ , and  $w_M$  are the values on  $[-1,1]$  and  $\xi_k$ ,  $\eta_L$ , and  $\eta_M$  are given by equations (17).

#### E. Temperature Distribution

One of the assumptions made in developing the equations herein was that the material properties are not functions of temperature. This assumption was made only to conserve computer time. The assumption does, of course, restrict the maximum temperatures to around three hundred degrees Fahrenheit for aluminum. This range of temperature is, however, more than sufficient to demonstrate the validity of the theory before large deflection effects become significant.

The program can handle either an analytical temperature "surface" or an experimentally measured temperature distribution. The analytical temperature distribution is specified in the form of a polynomial in the independent space variables as shown in Fig. 2. The only requirement for the measured temperature is that the measurements be made at a sufficient number of points to accurately define the temperature distribution.

The magnitude of the temperature distribution can be changed by inputting a series of  $\Delta T$ 's. In this case, since  $T_{ref}$  is used as 1.0, the  $\Delta T$ 's are input as the actual value of the temperature desired (in degrees Fahrenheit or Centigrade depending on the system of units used).

Any experimentally measured temperature distribution may be input. The values of the temperature at the integration points are calculated by a two-dimensional, quadratic interpolation subroutine. The temperatures are input on a rectangular grid. The points are evenly spaced in the  $\xi$  and  $\eta$  directions although the respective spacings need not be equal (i.e. the elements of the grid need not be square). A sample of the grid and an explanation of the defining parameters is shown in Fig. 3.

KC (I)	=	number of the first horizontal line at the $I^{th}$ vertical line ( $I = 1, 2, \dots, NTX$ )
LC(I)	=	the number of the last horizontal line at the $I^{th}$ vertical line ( $I = 1, 2, \dots, NTX$ )
DTX	=	distance between vertical lines
DTY	=	distance between horizontal lines

(XT1, YT1) = coordinates of lower left hand point of the grid.

NPTS = (not input) is calculated internally.  
This is the total number of grid points

For the grid in Fig. 3,

NTX = 7  
KC(I) = 1, 2, 2, 3, 3, 4, 4  
LK(I) = 14, 14, 13, 13, 12, 12, 11  
DTX = .142  
DTY = .13  
(XT1, YT1) = (0.0, - .8)

The temperatures at the grid points are input from bottom-to-top for each vertical line starting from the left side. This sequence is shown by the circled numbers in Fig. 3.

Interpolation will be attempted at any point within DTX and/or DTY of one of the grid points. This is a modification of a program contained in Ref. 18, in which a complete description is given.

The variable called TREF in the program is not actually used anywhere in the calculations. It is simply used as additional information to be output. Thus there are two ways of inputting the temperature distributions and incrementing the  $\Delta T$  values.

The first method is to simply input the actual magnitudes of the temperatures on the plate. In this case the values of  $\Delta T$  will be of 0 (1). At  $\Delta T=1^\circ$ , then, the eigenvalues calculated will give the frequencies of the plate for the input temperature distribution.

If desired, the temperature distribution may be normalized with respect to the temperature, TREF, at some reference point on the plate. In this case as with the analytical distribution, the  $\Delta T$ 's are input as the actual value of the temperature desired at the reference point.



#### F. Thickness Distribution

The thickness distribution  $h/h_p$ , is symmetric about the  $\xi$ - $\eta$  plane and is described by two polynomials in  $\xi$  and  $\eta$ . One gives the distribution on surface 1 and the other on surface 2. These two surfaces are separated by a line from the origin of the coordinate system at the angle  $\beta$ . The value  $h_0$  (called T0 in the program) is the thickness at the origin.

#### G. Miscellaneous Comments

The eigenvalue routine used here is a double precision version of the subroutine "NROOT" from the IBM Scientific Subroutine Package. (Note that this requires a double precision version of the subroutine "EIGEN" from the same source). The subroutine "DMINV" and "DGMPRD" (no listing given) are used directly from that source.

Extensive use was made of the disk storage available in writing the program. This reduced the core storage requirements to around 250,000 bytes on the IBM 370-165 computer used. Although the execution time for the program using all core storage would be about one-third of that using disk storage, the program would be limited to only thirty deflection and stress function terms and ten quadrature points. Also, core storage was a premium at the time of writing because of the large amount of business done by the Computer Center at The Ohio State University.

It should also be noted that for the coordinate system used some plates will be symmetric about the x-axis. In these cases the even and odd terms in the assumed solution uncouple. Thus, the deflection function may be separated into one function containing only even terms in  $\eta$  and one containing only odd terms in  $\eta$ . Each of these functions can then be input separately to give all even modes or all odd modes respectively. The same comments also apply to the stress function.

### IV. RESULTS

To compare to results in the literature, a conversion from the notation used in most other sources to that used herein is necessary. As long as the results are presented in non-dimensional form, only ratios of the material properties are required. Thus let,

$$E_{11}/E_{22} = D_x/D_y$$

$$G_{12}/E_{22} = D_k/D_y$$

$$E_{12}/E_{22} = D_{xy}/D_y - 2 D_k/D_y \quad .$$

All the data presented here will be converted using these relations to the notation previously described.

All of the computations presented in this section were made using a 36 mixed term deflection function,

$$w(\xi, \eta) = g(\xi, \eta) \sum_{i=0}^5 \sum_{j=0}^5 h_{ij} \xi^i \eta^j \quad .$$

Thus, the first thirty-six modes and frequencies were calculated. The runs took an average of three minutes (Central Processing Unit) time on the IBM 370-165.

No effort was made to optimize the program, the purpose being to obtain consistently good results for any planform shape with any boundary condition. e.g., acceptable results can be obtained for the torsion mode of a rectangular cantilever plate with only three terms in the deflection function. However, this number of terms is completely inadequate for any other mode of the thermally stressed cantilever plate and is inadequate for any mode of any other of the many plates investigated. Thus, the large number of terms in both the stress function and the displacement function may be much greater than required for some of the problems solved. This point is immaterial when the choice boils down to either obtaining an accurate quantitative answer in which one can have confidence or some answer that may only be in the "ball park."

#### A. Comparison of Orthotropic Results Without Thermal Loading

As was previously stated, the material published without thermal loading is voluminous. For the sake of brevity, only a few comparisons will be made.

Tables 1 and 2 are comparisons of calculated frequencies from the literature with those calculated by this program. It can be seen that the method under discussion here gives

excellent agreement with those frequencies. The expression for  $\lambda_1^2$  is given in the list of symbols and in Appendix A.

Table 3 gives a comparison with some experimental frequencies for plywood plates. Note that the experimental values are higher than the calculated frequencies. Because the method used here gives solutions which converge from above the exact solution, these errors are attributed to restraints inherent in the experimental approximation of the simply supported boundary conditions.

#### B. Comparison of Isotropic Results With Thermal Loading

As was stated previously, no experimental data were found in the literature on the effect of thermal stresses on orthotropic plates. Thus, a brief quantitative comparison is made with experimental data from Ref. 16 for the special case of the isotropic plate. The purpose is to show the agreement that was obtained for widely different cases. The isotropic plate elastic properties requires that,

$$E_{11} = E_{22} = E/(1-\nu^2)$$

$$E_{12} = \nu E/(1-\nu^2)$$

$$G_{12} = E/2(1+\nu)$$

As in Ref. 16, nominal values of the plate material properties used were,

$$E = 10^7 \text{ psi}$$

$$\nu = \frac{1}{3}$$

$$\alpha = 12.8 \times 10^{-6} / ^\circ\text{F}$$

Fig. 4 shows a comparison for the first five modes of a square cantilever plate.<sup>16</sup> It is interesting that the fourth and fifth mode frequency curves cross. A typical experimentally measured temperature distribution resulting from radiant lamp edge heating is shown in Fig. 5.

Fig. 6 presents an unsymmetrical trapezoidal cantilever plate that does not appear in Ref. 16. Only the first two modes were recorded for this plate. One of the temperature distributions measured on this plate is shown in

Fig. 7.

Because of its boundary conditions and the choice of coordinate system, the plate shown in Fig. 8 is also unsymmetrical. The stresses as well as the deflections are affected by the boundary conditions.<sup>16</sup> Here again, as in Fig. 4, two of the frequencies cross. The heating elements were centered over the diagonal from lower-left to upper-right giving a temperature distribution as shown in Fig. 9.

The frequencies of a plate with a single point clamped is shown in Fig. 10. The agreement with the four modes measured is seen to be good.

A plate with homogeneous pinned-free boundary conditions is shown in Fig. 11. Only the first two modes were recorded for this plate. The third calculated mode frequency is also shown. The temperature distribution shown in Fig. 5 is also typical of that used to calculate the frequencies for the plates in Figs. 10 and 11.

#### C. Comparison of Orthotropic Results With Thermal Loading

Figs. 12, 13, 14, 15, 16 and 17 constitute the results of a very brief study that indicates the large effect that orthotropy can have in the presence of thermal gradients. For the sake of brevity, only one boundary condition, the cantilever plate, and only one assumed temperature distribution,  $T = \Delta T |\eta|^3$ , is presented. Also, only the two lowest modes are presented although as many of the higher modes as could possibly be desired are available in the computer print-out. It should be noted that orthotropy does not change the characteristic shape of the response curves shown. However, because of the influence that the directional properties of the material can have on the stress field for a given temperature distribution, orthotropy can produce marked increases in the loss of effective stiffness for a given heating rate. In each figure the isotropic response curves are given for comparison purposes.

In Figs. 12 and 13, it can be seen that doubling the thermal expansion coefficient,  $\alpha_2$ , in the chordwise direction has little effect on the frequency for this temperature distribution because the chordwise stress component effect is small for this plate and remains small even when  $\alpha_2$  is doubled. However, when the longitudinal coefficient,  $\alpha_1$ , is doubled, the longitudinal stress component is essentially doubled and the frequency is seen to decrease at a markedly higher rate. Thermal buckling is reached at a temperature only one half as great as before. This means that, for a given heating rate, if an isotropic plate buckles within

thirty seconds, an orthotropic plate with this ratio of thermal coefficients would buckle in only fifteen seconds.

In Fig. 14 it can be observed that doubling the chordwise modulus of elasticity,  $E_{22}$ , actually produces a decrease in the rate of frequency decay and an increase in  $\Delta T_{cr}$  for the bending mode. That the various modes will not behave in the same way for a given material property change is shown in Fig. 15 where the doubled chordwise modulus produces the opposite effect on the torsion mode. Doubling the spanwise modulus,  $E_{11}$ , causes a higher rate of frequency decay with correspondingly lower buckling temperatures in both modes. It should be noted, however, that, although large, doubling the longitudinal modulus of elasticity does not have the extreme effect that is caused by doubling the longitudinal conductivity coefficient.

Fig. 16 shows that the decrease in the rate of decay of the bending mode, achieved by doubling  $E_{22}$  in Fig. 14, is more than offset by the increase caused by doubling  $\alpha_1$  in Fig. 12. Both Figs. 16 and 17 show appreciable destabilizing effects when both the moduli and expansion coefficient are changed. It is recognized that the change in properties as used is large, but the effects are also large. Using the computer program presented herein, almost unlimited parametric studies may be made to determine trends or, more efficiently, calculations may be made to obtain answers for specific cases once a problem has been defined.

#### D. Effect of Stress Distribution

An interesting by-product of this program is an accurate calculation of the thermal stress distribution. The effect of the stress distribution function used in calculating the vibration frequency is shown in Fig. 15. The boundary conditions require a stress function containing both odd and even terms in  $\eta$ . Thus, the stress distribution calculated, using only even terms, does not give the correct stress distribution and hence, does not give the correct frequency. However, when the correct mixed terms were used, the correct stress distribution existing in the plate resulted and consequently, the calculated frequencies agreed more closely with the experimental values. The 36 mixed term results show that the analytically predicted stress distribution had effectively converged when 24 mixed terms were used.

The three sets of frequencies calculated were compared to the experimental response of the plate shown in Fig. 8. The results show that correct stresses are in fact necessary in order to obtain the correct frequencies.

## V. CONCLUDING REMARKS

The computer program presented herein is, in effect, a general solution to the problem of the linear vibration of thermally stressed trapezoidal plates. The theory has been verified experimentally for thermally stressed isotropic plates and has been found to agree favorably with analytical data found in the literature for orthotropic plates with no thermal loading. It appears that accurate results can be obtained by the methods herein described for almost any boundary conditions of practical importance. The solutions, based on linear theory, do not hold as the buckling region is approached because of the non-linear effects of large deflections.

Experience in using the program shows that the number of terms required in the assumed solutions increases with the complexity of the geometry. However, consistently accurate results are obtained using a 30-36 term displacement function and a 24-30 term stress function. Accurate integration can be obtained for this number of terms using ten quadrature points in the  $\xi$ -direction (twenty points in the  $\eta$ -direction).

A very extensive experimental program with orthotropic plates would be required to verify all the facets of the program as presented. However, the data comparisons shown combined with many cases of experimental isotropic plate data not presented herein gives the writers great confidence in the analytical results.

The generality of this program should not be overlooked. Its extension to obtain the accurate solution of the flutter of thermally stressed plates and panels can be readily made. A natural extension of this work would be to examine, without the assumption of mode identity, the large deflection effects observable as heating progresses. A recently developed method of solving large sets of non-linear equations shows great promise in the area of large deflections which has yet to be effectively investigated.

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TABLE 1  
PF-CC-PF-CC

AR	$E_{11}$	$E_{12}$	$E_{22}$	$G_{12}$	$\lambda_1$		Ref. No.
					Calculated By Program	Reference Values	
1.0	1.0	1.0	1.0	0.0	28.93	29.29	8
1.0	2.0	1.0	1.0	0.0	21.6	21.82	8
1.0	1.0	3.0	1.0	0.0	36.1	36.5	8
1.0	3.0	3.0	1.0	0.0	22.35	22.56	8
1.0	6.0	2.0	1.0	0.0	16.12	16.13	8
1.0	1.0	1.0	2.0	0.0	50.8	51.5	8
1.0	1.0	3.0	9.0	0.0	73.0	74.0	8
2.0	3.117	0.12	1.0	0.264	53.6	53.7	3

TABLE 2

PF-FF-PF-CF AR = 2.0

 $E_{11} = 3.177$ ,  $E_{12} = 0.12$ ,  $E_{22} = 1.0$ ,  $G_{12} = 0.264$ 

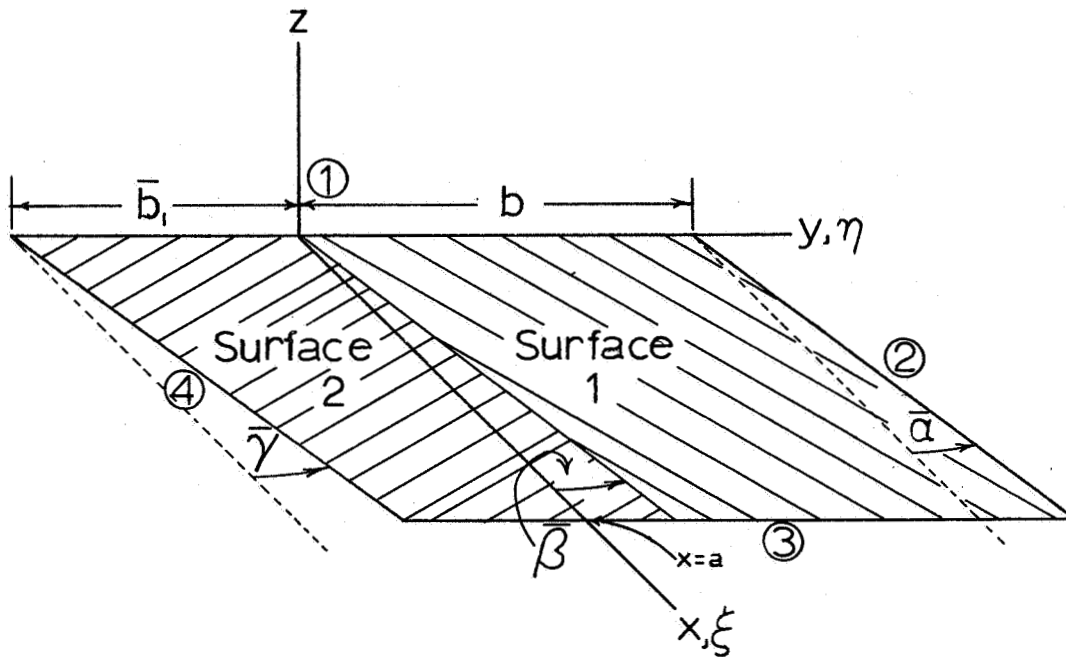
Mode No.	$\lambda_1$	
	Calculated By Program	REF 21 Values
1	14.75	14.75
3	55.35	55.4
5	93.1	91.6
6	121.4	120.0
7	144.4	144.1
12	256.1	249.0
13	295.1	278.2

TABLE 3

PF-PF-PF-PF AR = 1.0, a = 45.8 cm.

$E_{11} \times 10^{-10}$ Dynes cm <sup>2</sup>	$E_{12} \times 10^{-10}$ Dynes cm <sup>2</sup>	$E_{22} \times 10^{-10}$ Dynes cm <sup>2</sup>	$G_{12} \times 10^{-10}$ Dynes cm <sup>2</sup>	$\rho$ gm. cc.	h cm.	$f_1$ (cps)	
						Calc.	REF. 2
6.9	0.17	0.17	0.30	0.33	0.291	32.06	35
7.4	0.17	0.05	0.34	0.39	0.323	33.8	39
7.9	0.19	0.19	0.45	0.399	0.310	34.0	34
13.3	0.33	0.55	0.85	0.67	0.309	34.5	41

# General Planform Parameters and Non-Dimensionalization



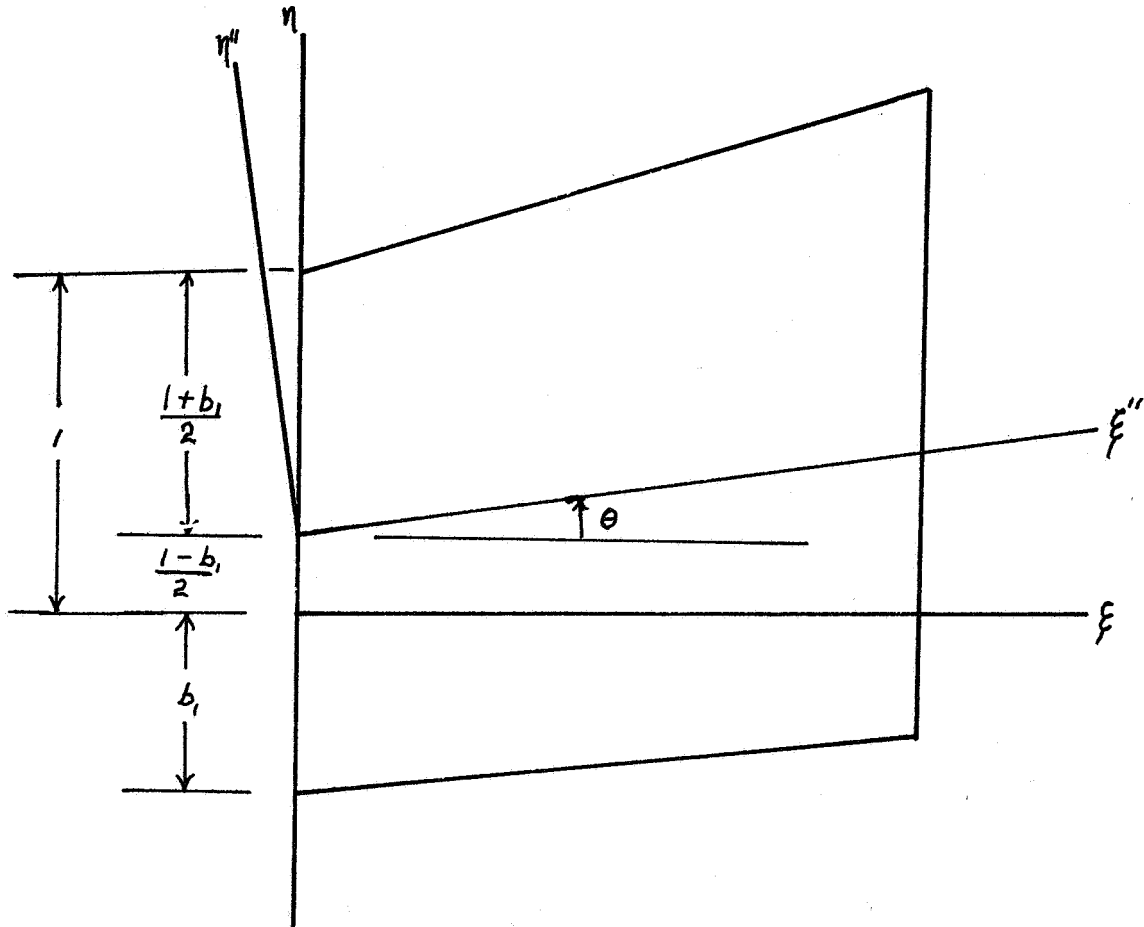
$$\xi = \frac{x}{a} \quad \eta = \frac{y}{b}$$

$$\alpha = \frac{a}{b} \tan \bar{\alpha} \quad \beta = \frac{a}{b} \tan \bar{\beta} \quad \gamma = \frac{a}{b} \tan \bar{\gamma}$$

$$b_1 = \frac{\bar{b}_1}{b} \quad AR = \frac{a/b}{1 + b_1 - (\frac{\gamma - \alpha}{2})}$$

Fig. 1

# Analytical Temperature Distribution



$$T(\xi, \eta) = \sum_{I=1}^{NTEMP} TEM(I)(\eta'')^{NTEM(I)}(\xi)^{NTEMX(I)}$$

$$\eta = \frac{|\eta''|}{\frac{l+b_1}{2} \cos \theta}$$

Fig. 2

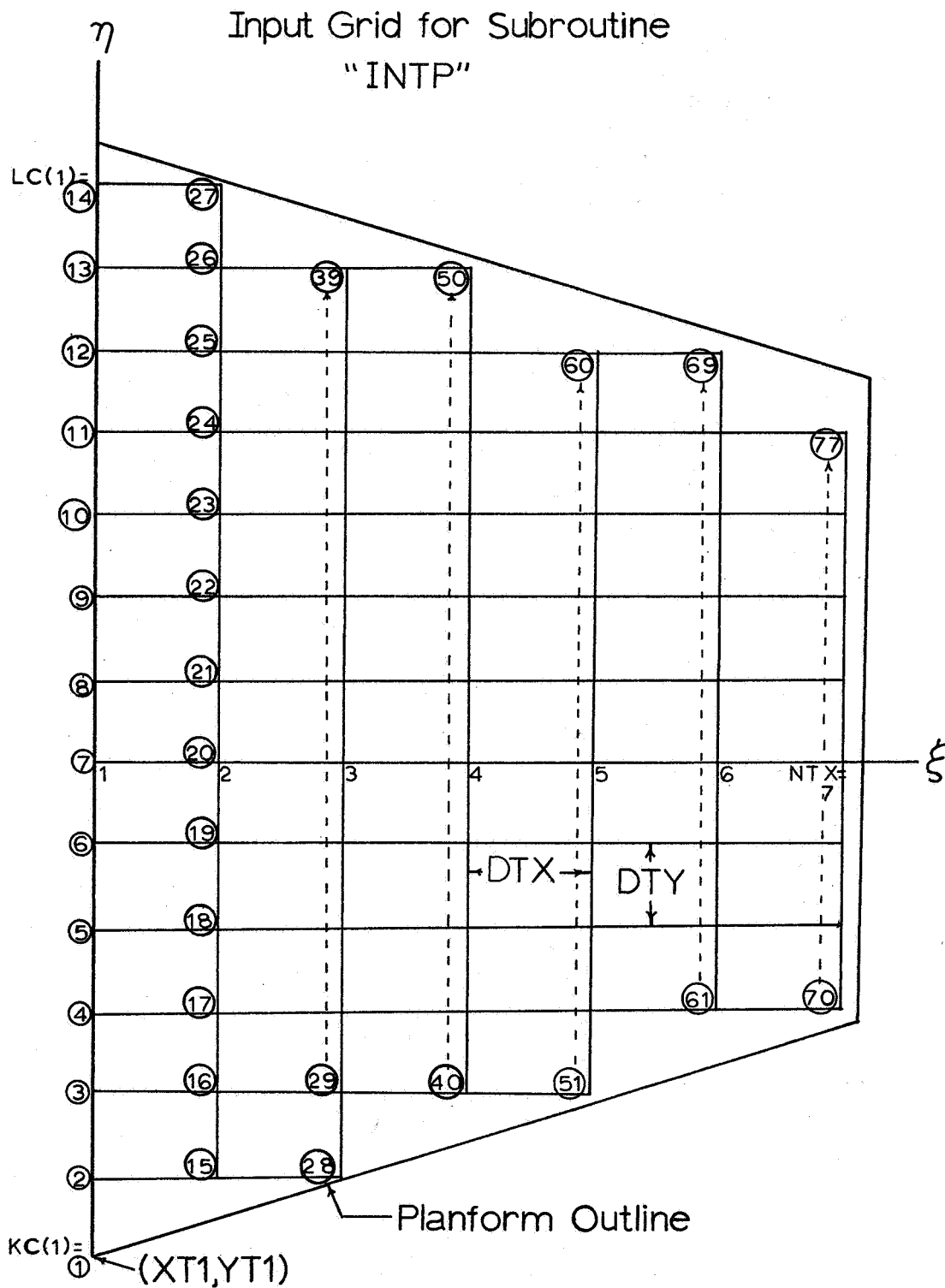
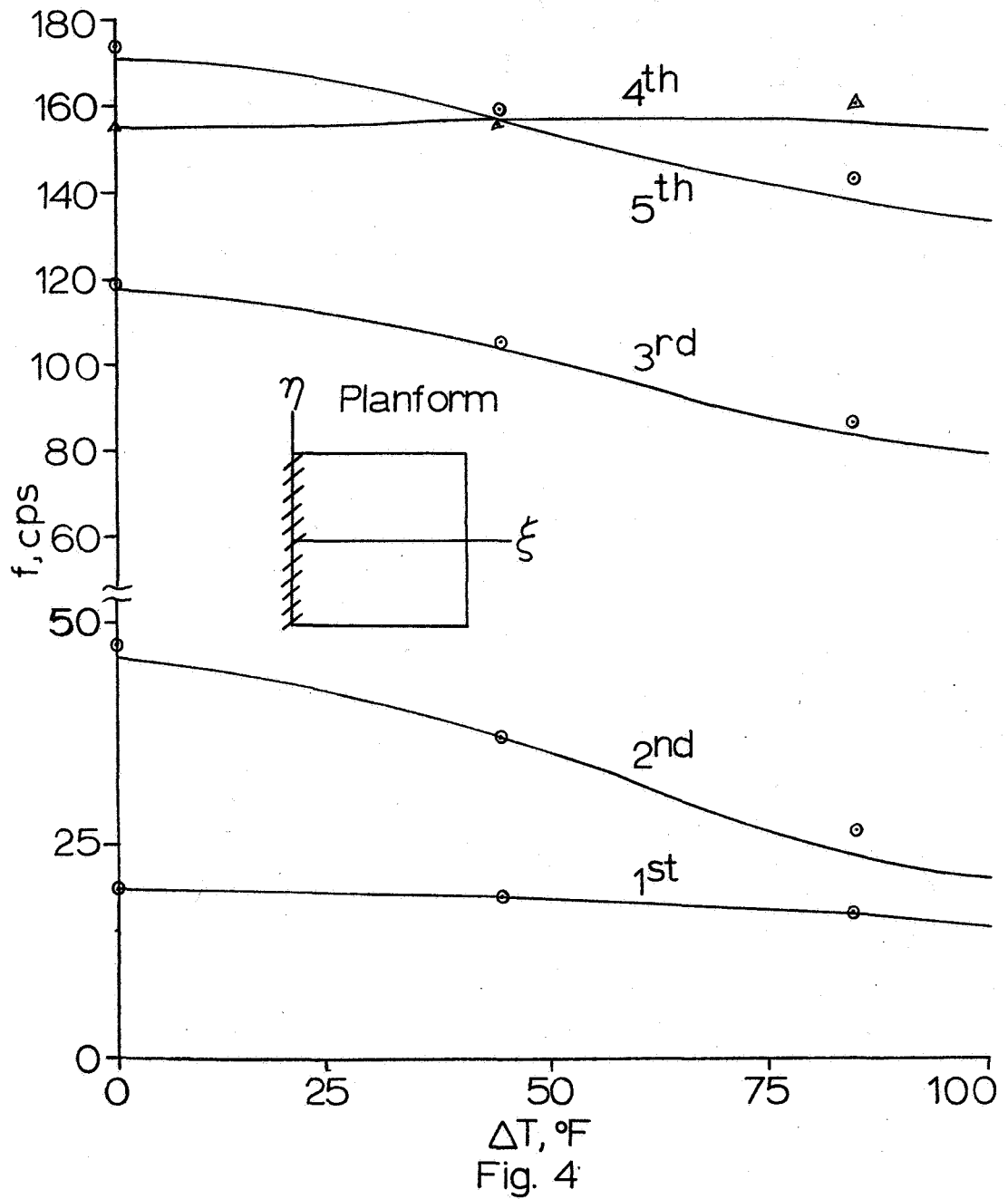


Fig. 3

Comparison of Analytical and  
Experimental Response of a  
CC-FF-FF-FF Plate  
AR=1.0 , a=18" , h=3/16"

— measured  
• calculated



Typical Temperature  
Distribution used in  
Fig. 4,10,11

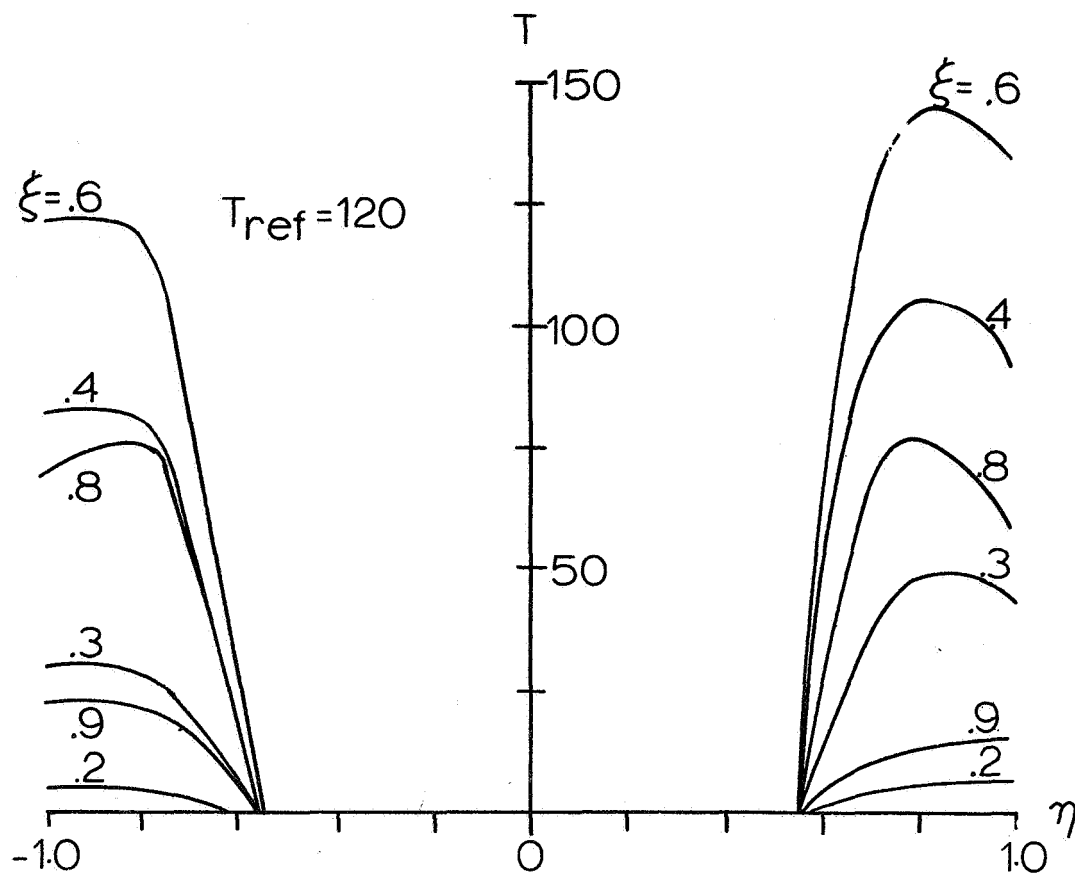


Fig. 5

Comparison of Analytical and  
Experimental Response of a  
CC-FF-FF-FF Plate

$$\alpha = -0.6 \quad \beta = 0.0 \quad \gamma = 0.0$$

$$AR = 5/3, \quad a = 20", \quad h = 1/4"$$

— measured

• calculated

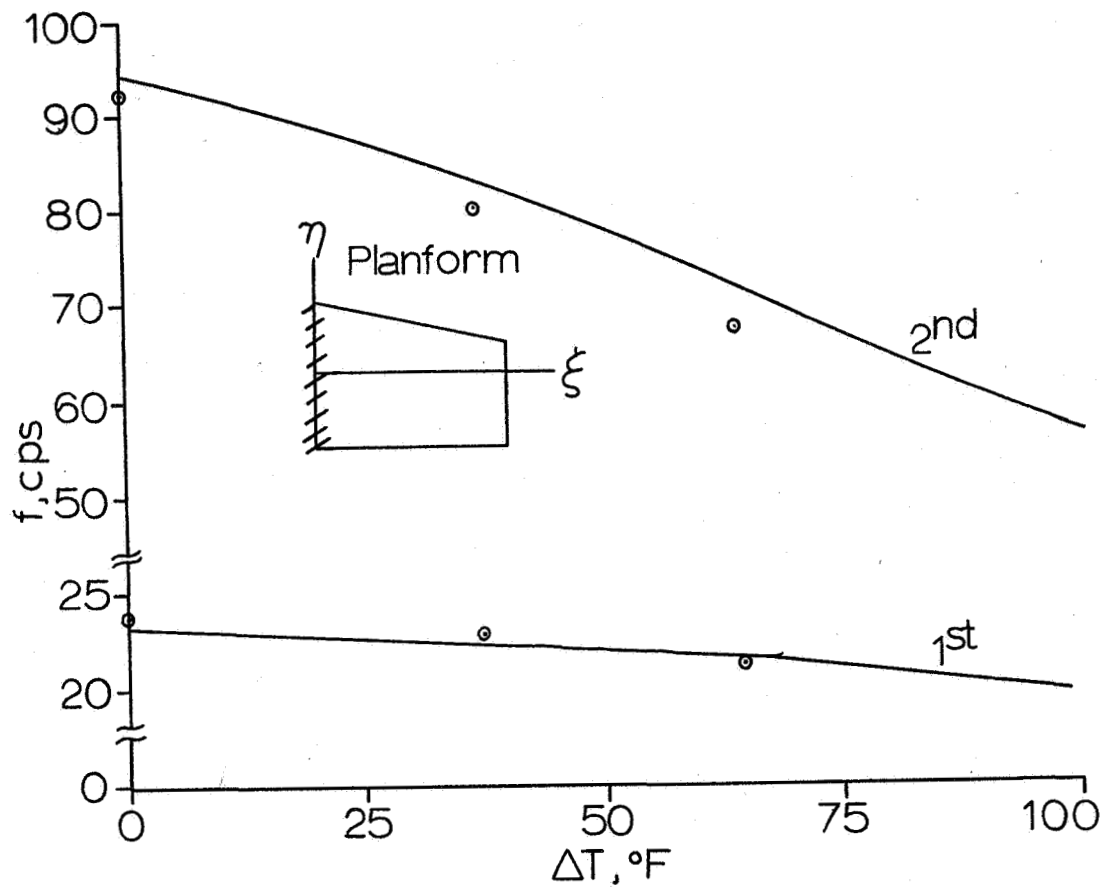


Fig. 6



Temperature Distribution for Fig. 6

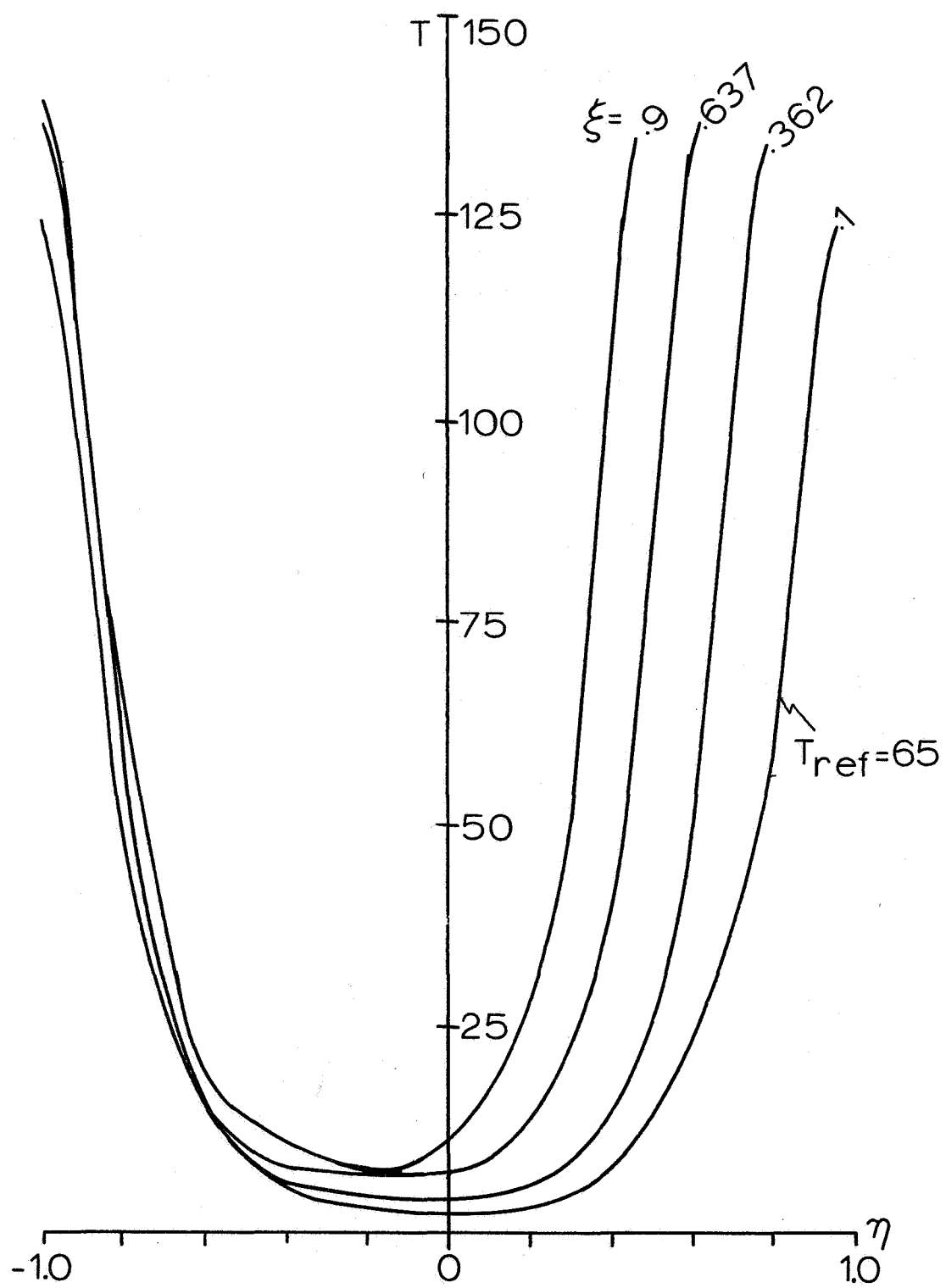


Fig. 7

Comparison of Analytical and  
Experimental Response of a  
CC-FF-FF-CC Plate  
AR=1.0,  $a=18''$ ,  $h=3/16''$

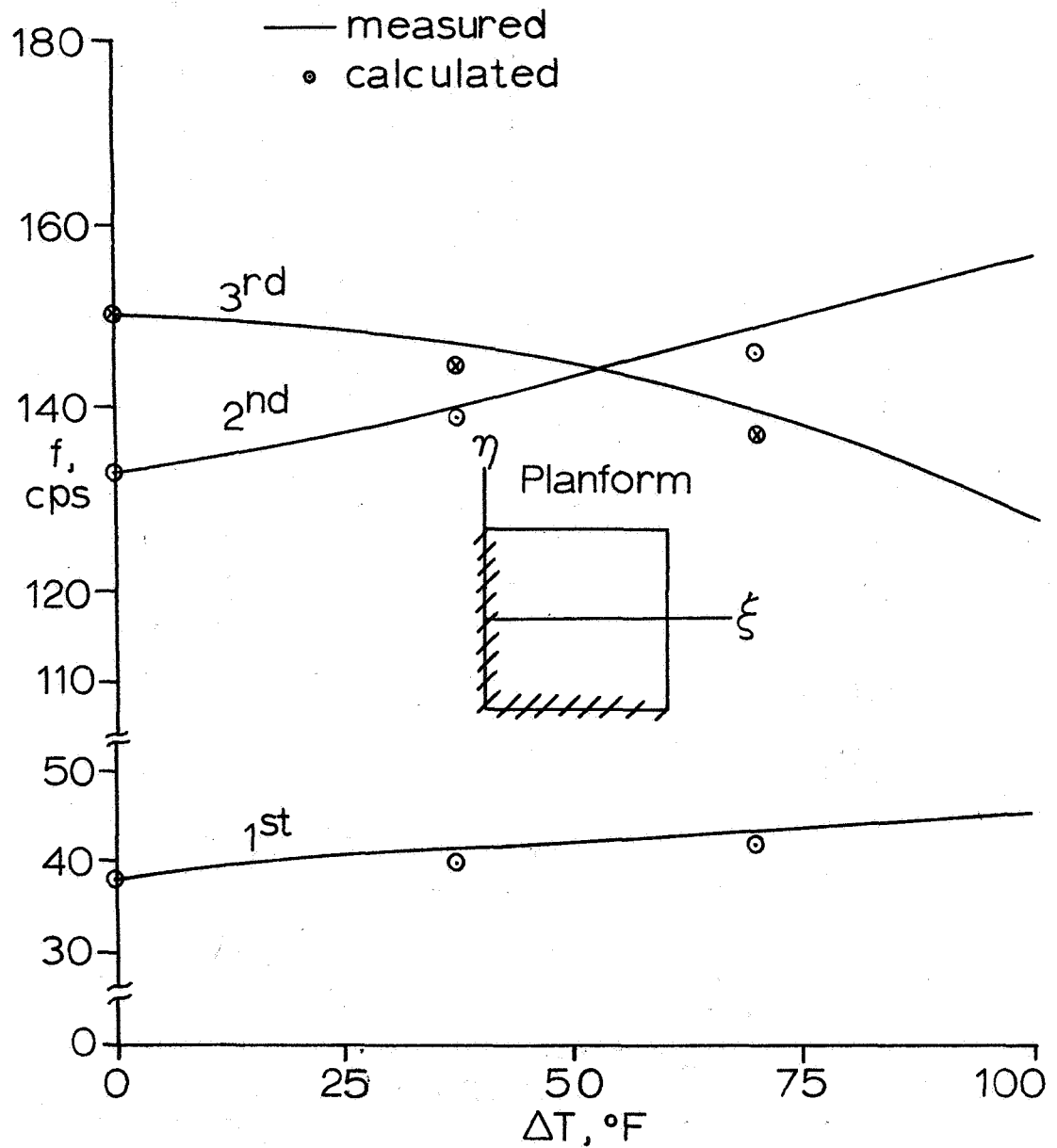


Fig. 8

Typical Temperature  
Distribution used in  
Fig. 8

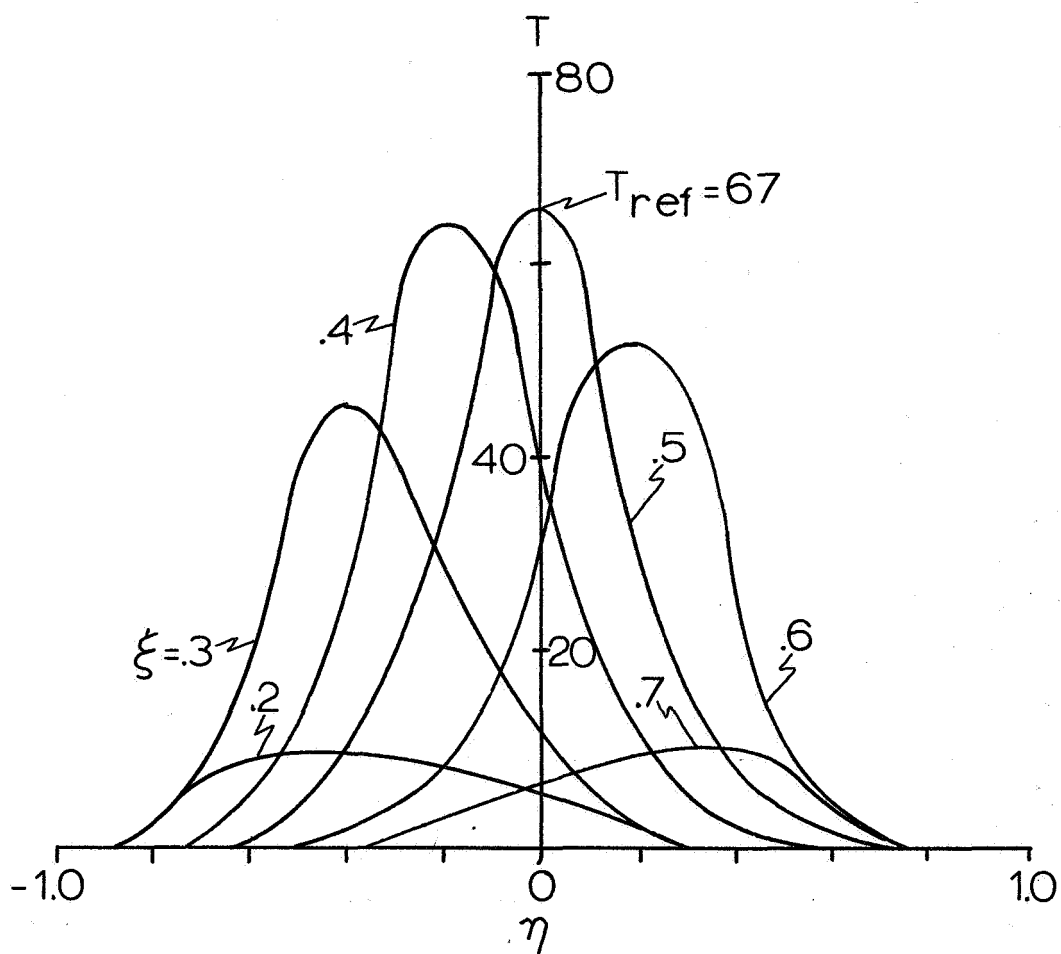


Fig. 9

Comparison of Analytical and  
Experimental Response of a  
Plate Clamped at (0,0)

AR=1.0 , a=18" , h=3/16"

— measured

• calculated

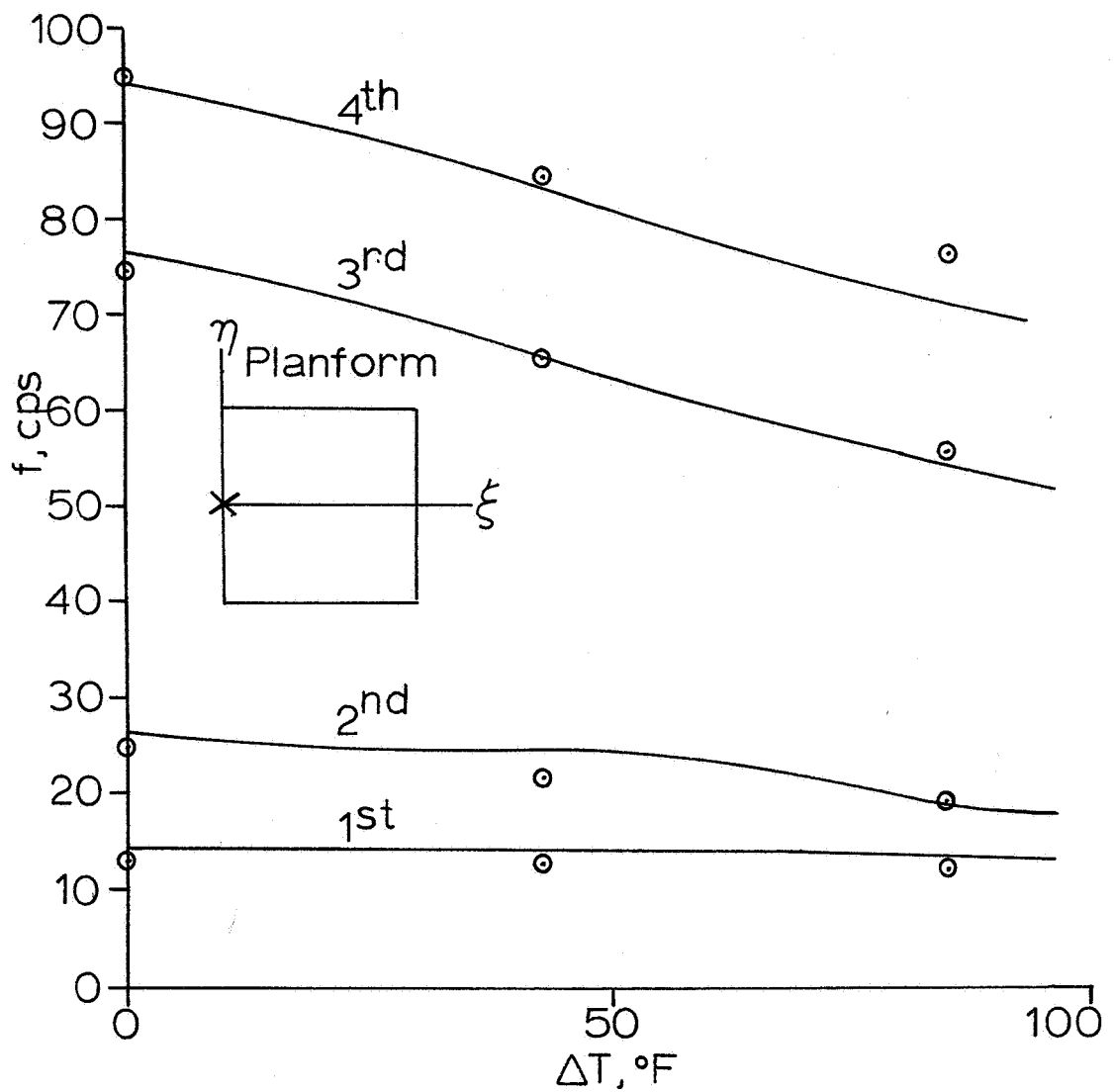


Fig. 10

Comparison of Analytical and  
Experimental Response of a  
PF-PF-PF-PF Plate  
AR=1.0 , a=18" , h=3/16"

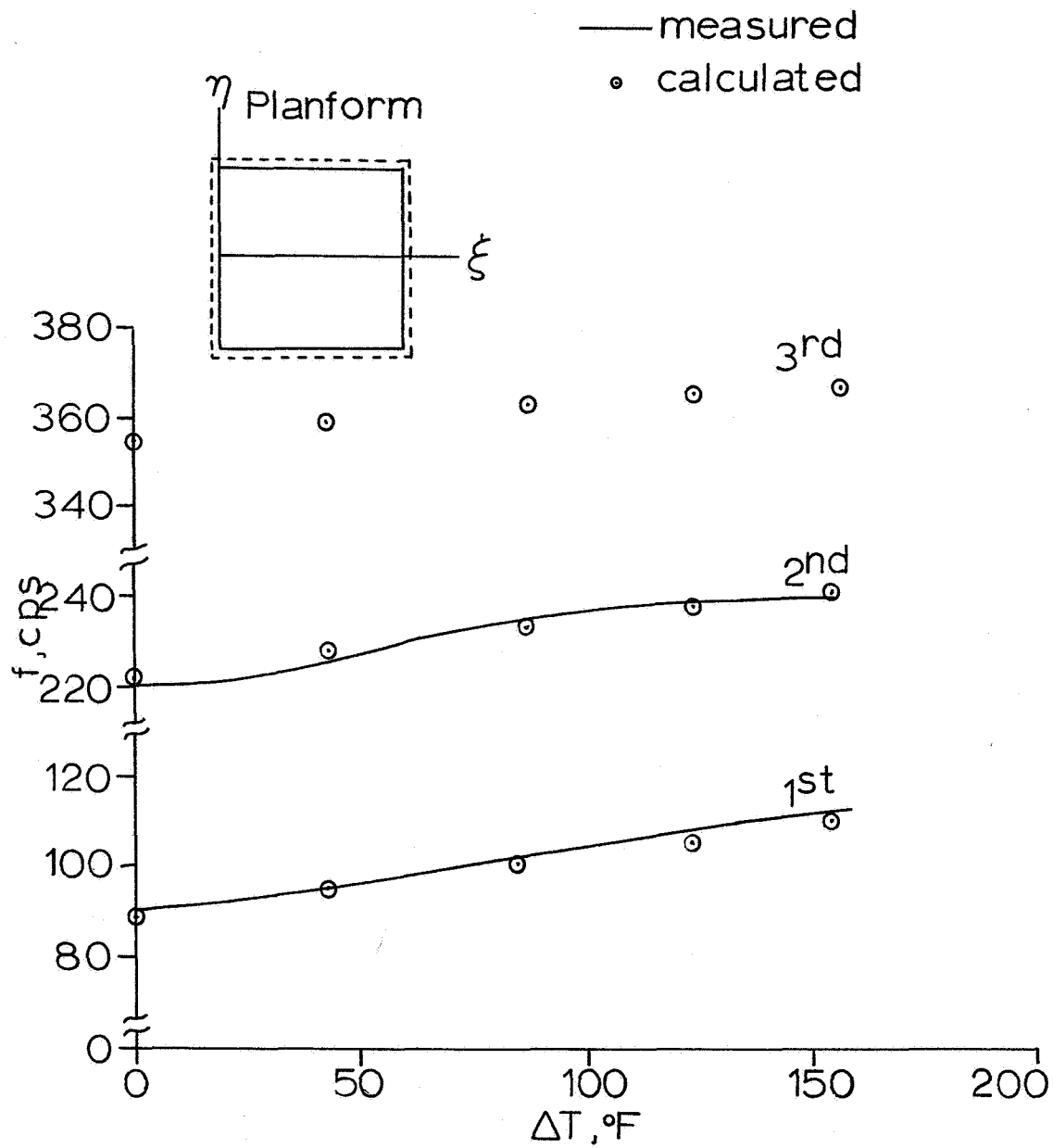


Fig. 11

# Effect of $\alpha$ on Plate Vibration

$$T = \Delta T |\eta|^3$$

Constant Thickness

$$AR = 5/3$$

1<sup>st</sup> mode

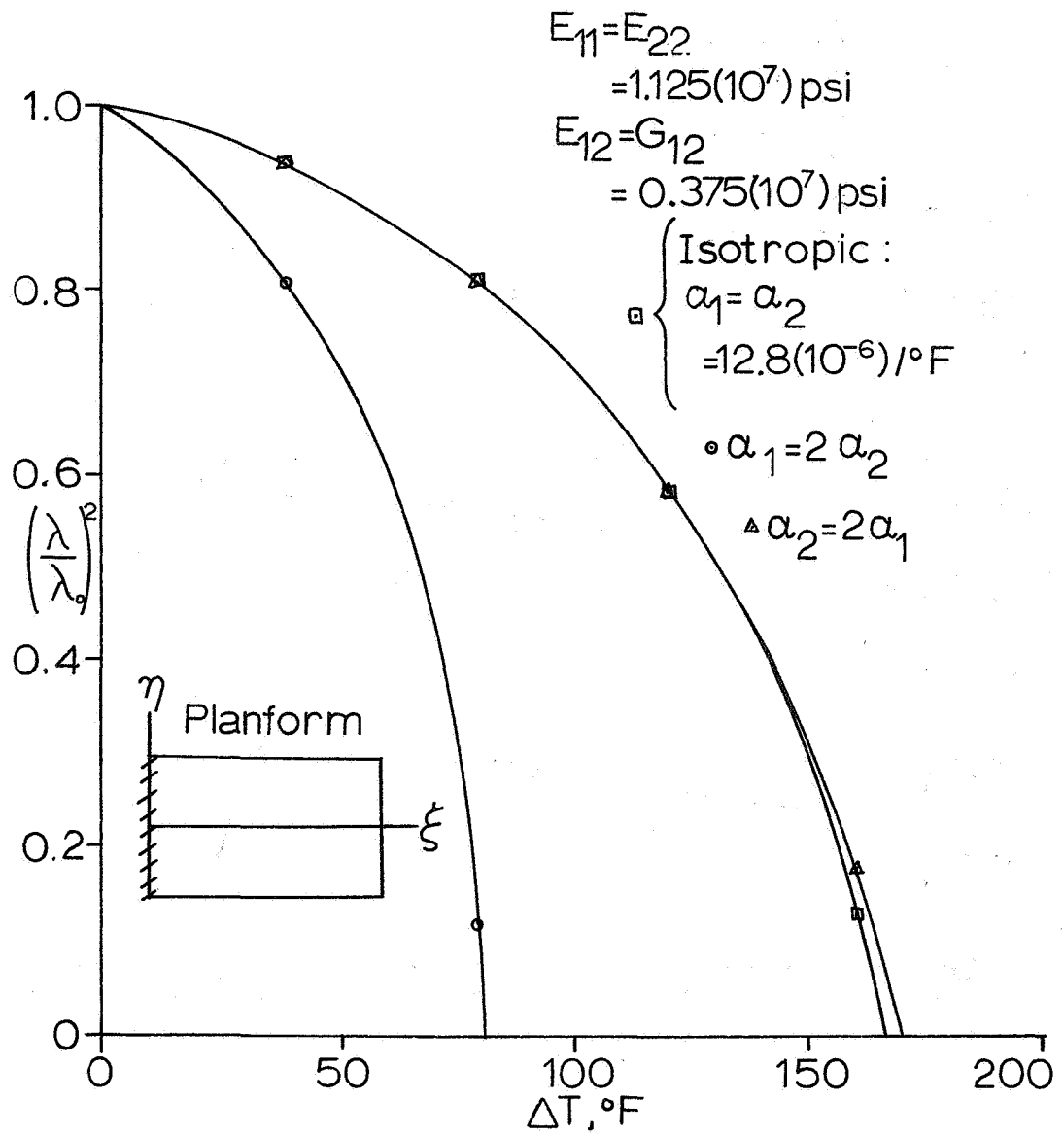


Fig. 12

Effect of  $\alpha$  on Plate  
Vibration  
2<sup>nd</sup> mode  
(See Fig.12 for notation)

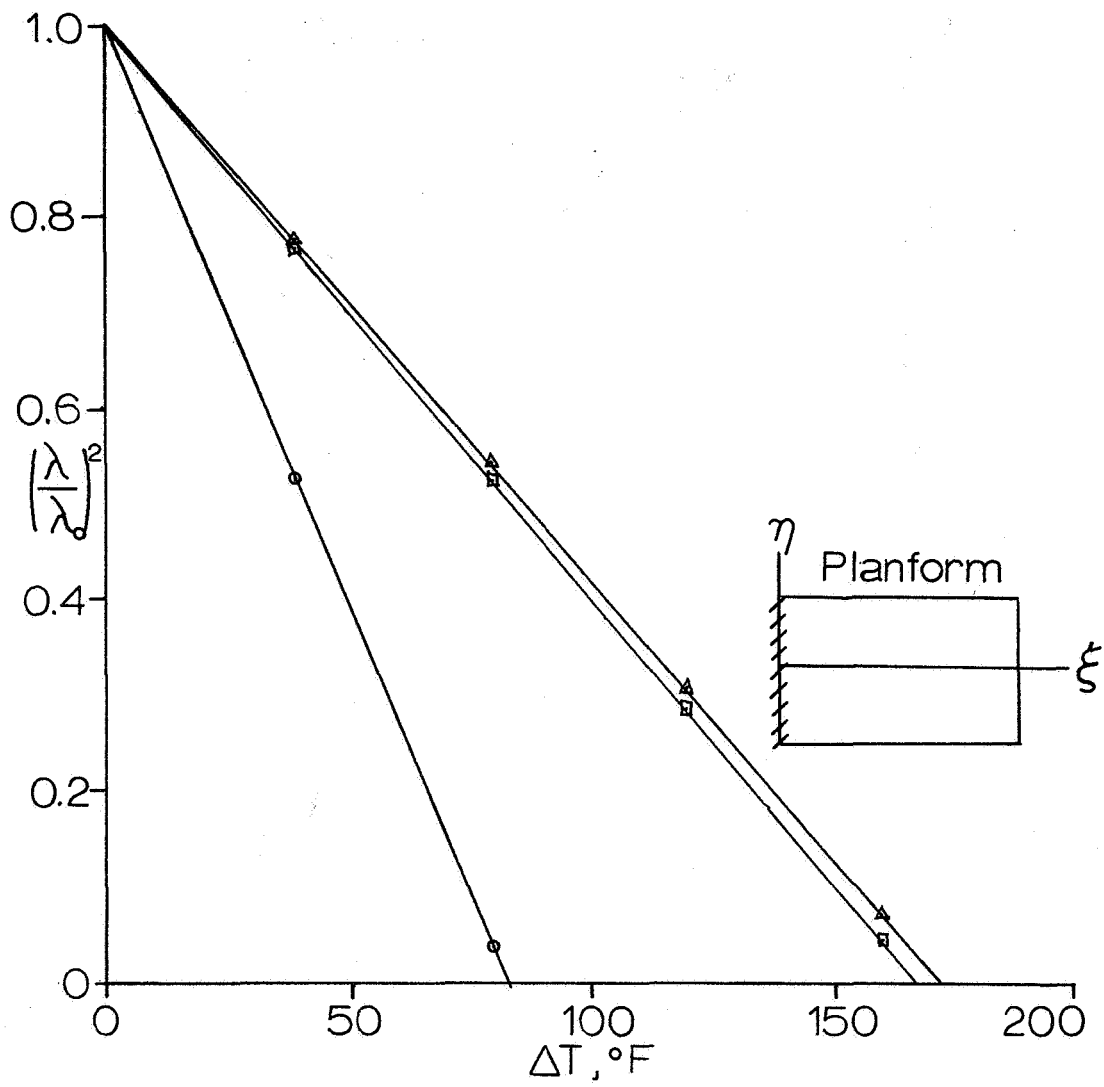


Fig.13

# Effect of 'E' on Plate Vibration

AR= 5/3

Constant Thickness

$$T = \Delta T \eta l^3$$

1<sup>st</sup> mode

$$E_{12} = G_{12} = 0.375(10^7) \text{ psi}$$

$$\alpha_1 = \alpha_2 = 12.8(10^{-9}) / ^\circ\text{F}$$

$$\square \left[ \begin{array}{l} \text{Isotropic:} \\ E_{11} = E_{22} = 1.125(10^7) \text{ psi} \end{array} \right.$$

$$\circ E_{22} = 2E_{11}$$

$$\triangle E_{11} = 2E_{22}$$

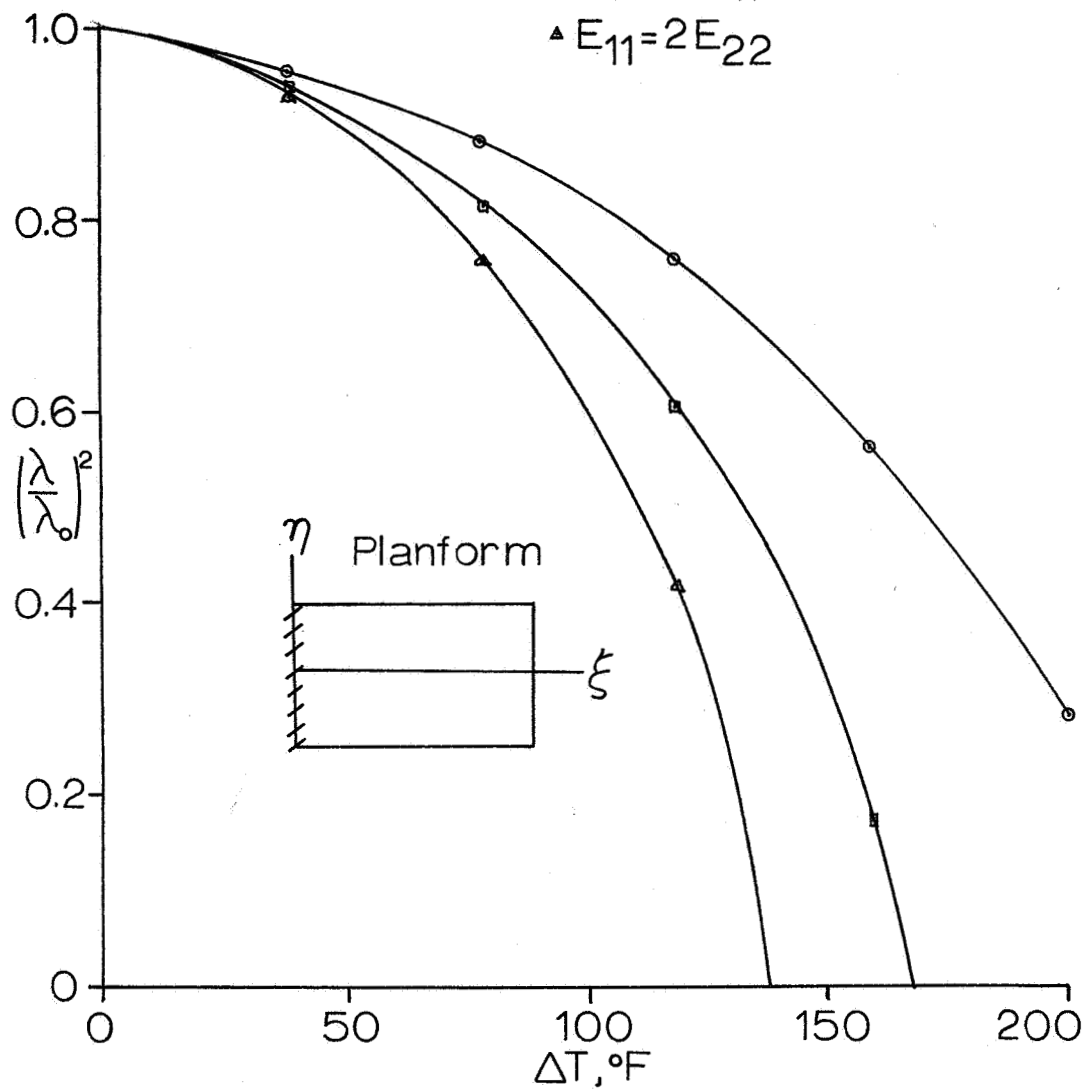
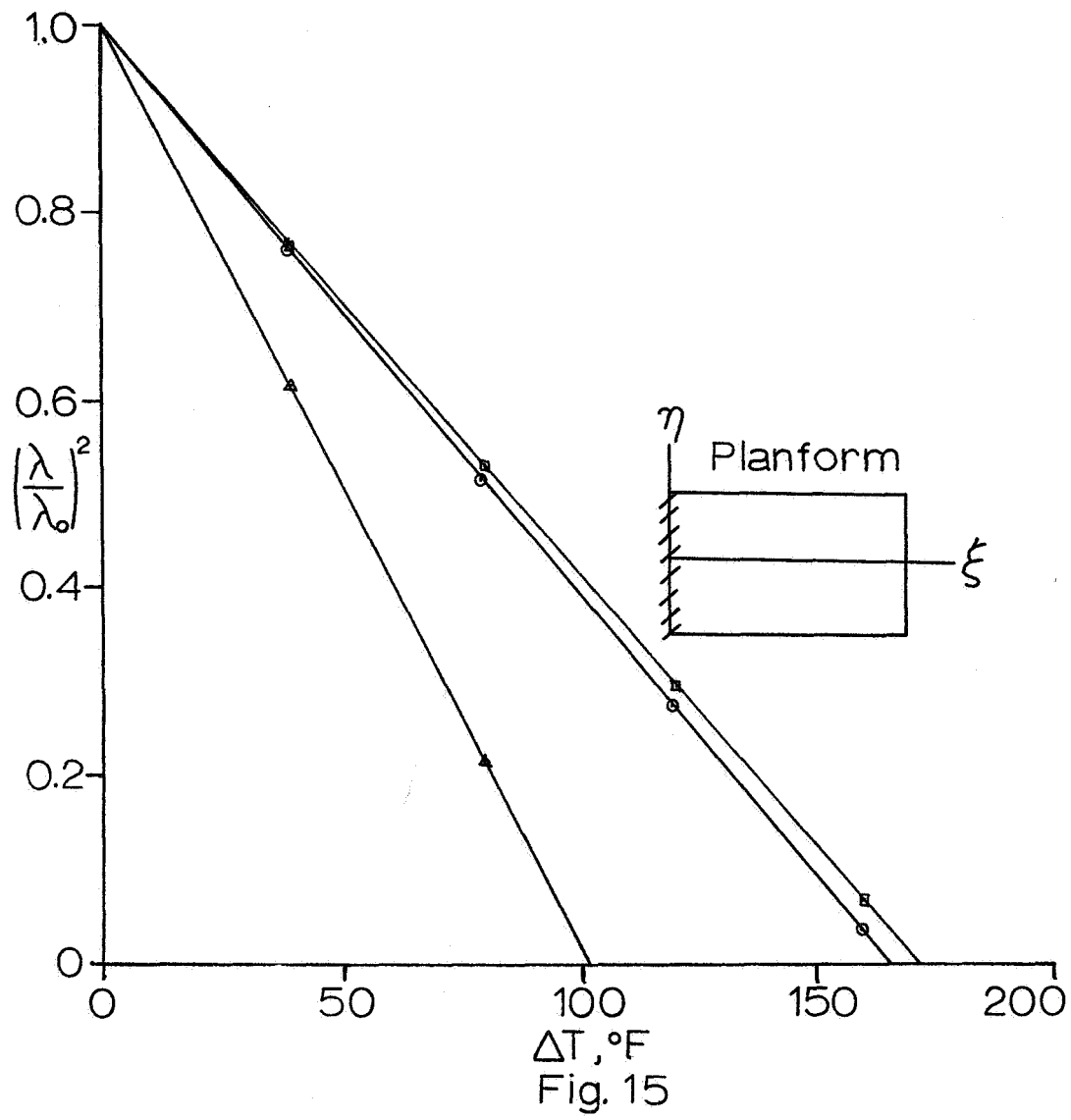


Fig. 14



Effect of 'E' on Plate  
Vibration  
2<sup>nd</sup> mode  
(See Fig 14 for notation)



# Combined Effect of 'E' and $\alpha$ on Plate Vibration

$$AR = 5/3$$

Constant Thickness

$$T = \Delta T |\eta|^3$$

1<sup>st</sup> mode

$$\begin{cases} E_{12} = G_{12} = 0.375(10^7) \text{ psi} \\ \text{Isotropic:} \\ E_{11} = E_{22} = 1.125(10^7) \text{ psi} \\ \alpha_1 = \alpha_2 = 12.8(10^{-6}) / ^\circ\text{F} \\ \bullet E_{11} = 2E_{22}, \alpha_2 = 2\alpha_1 \\ \blacktriangle E_{22} = 2E_{11}, \alpha_1 = 2\alpha_2 \end{cases}$$

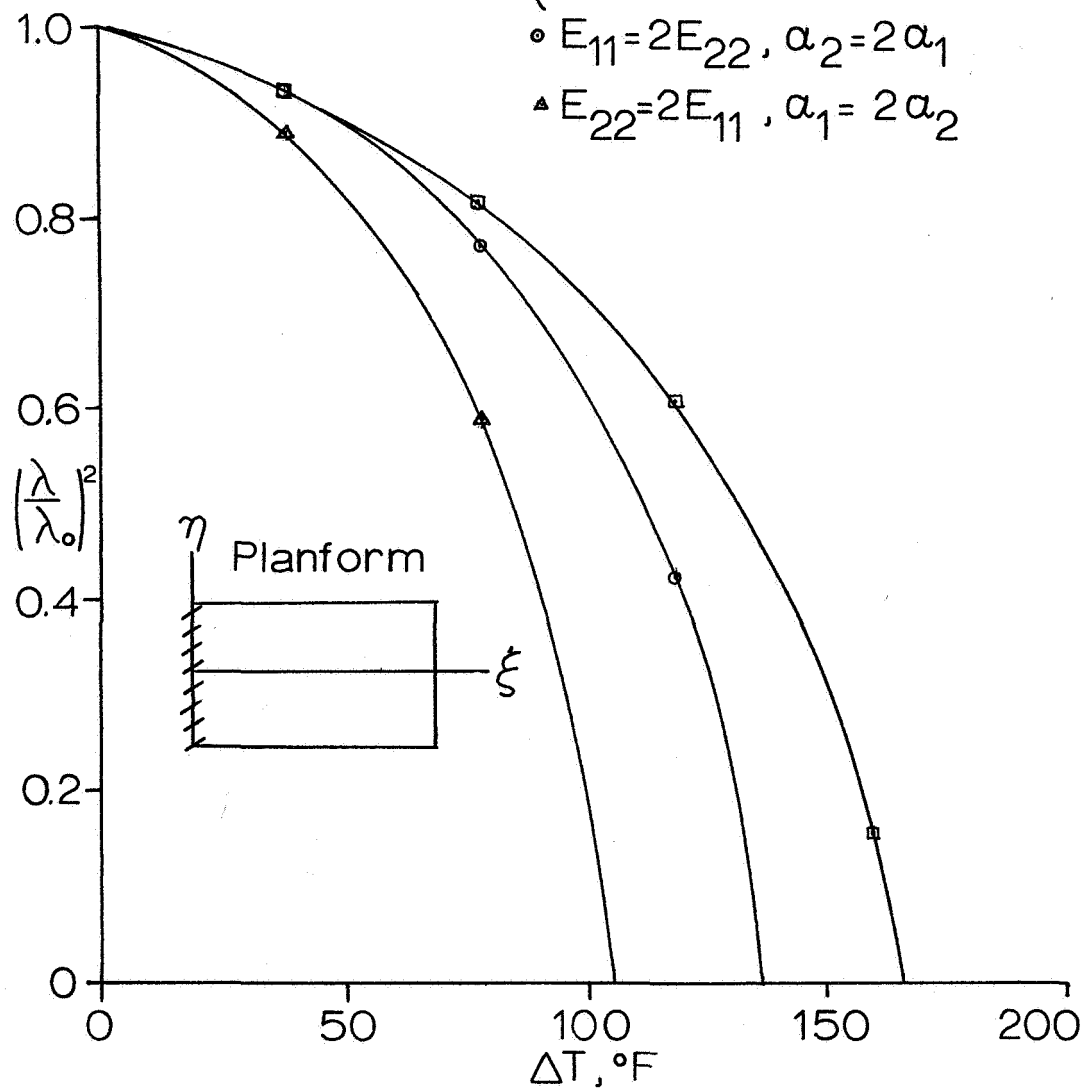
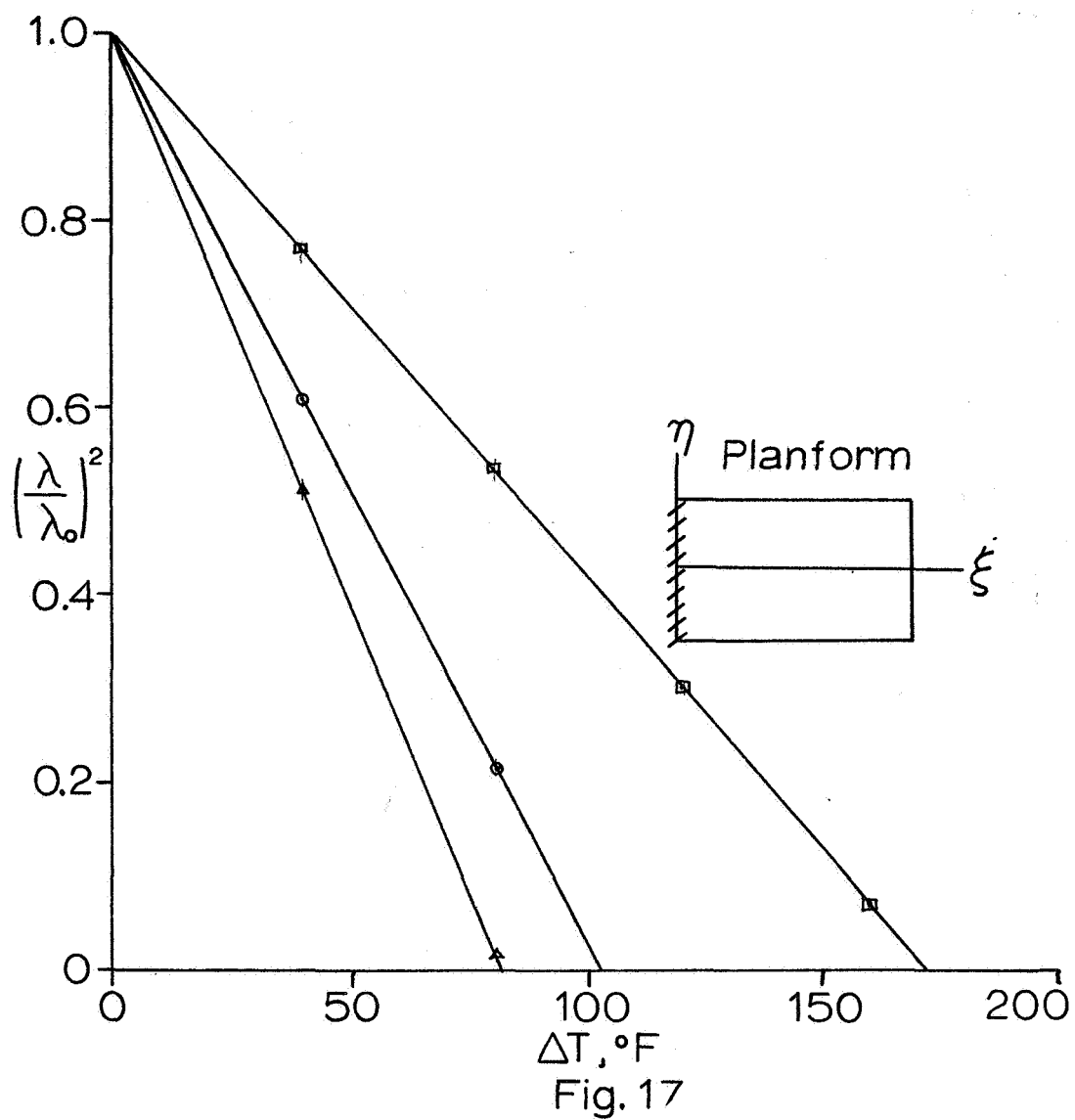


Fig. 16

Combined Effect of 'E' and  $\alpha$   
on Plate Vibration  
2<sup>nd</sup> mode  
( See Fig.16 for notation)



# Effect of Stress Function on Plate Vibration

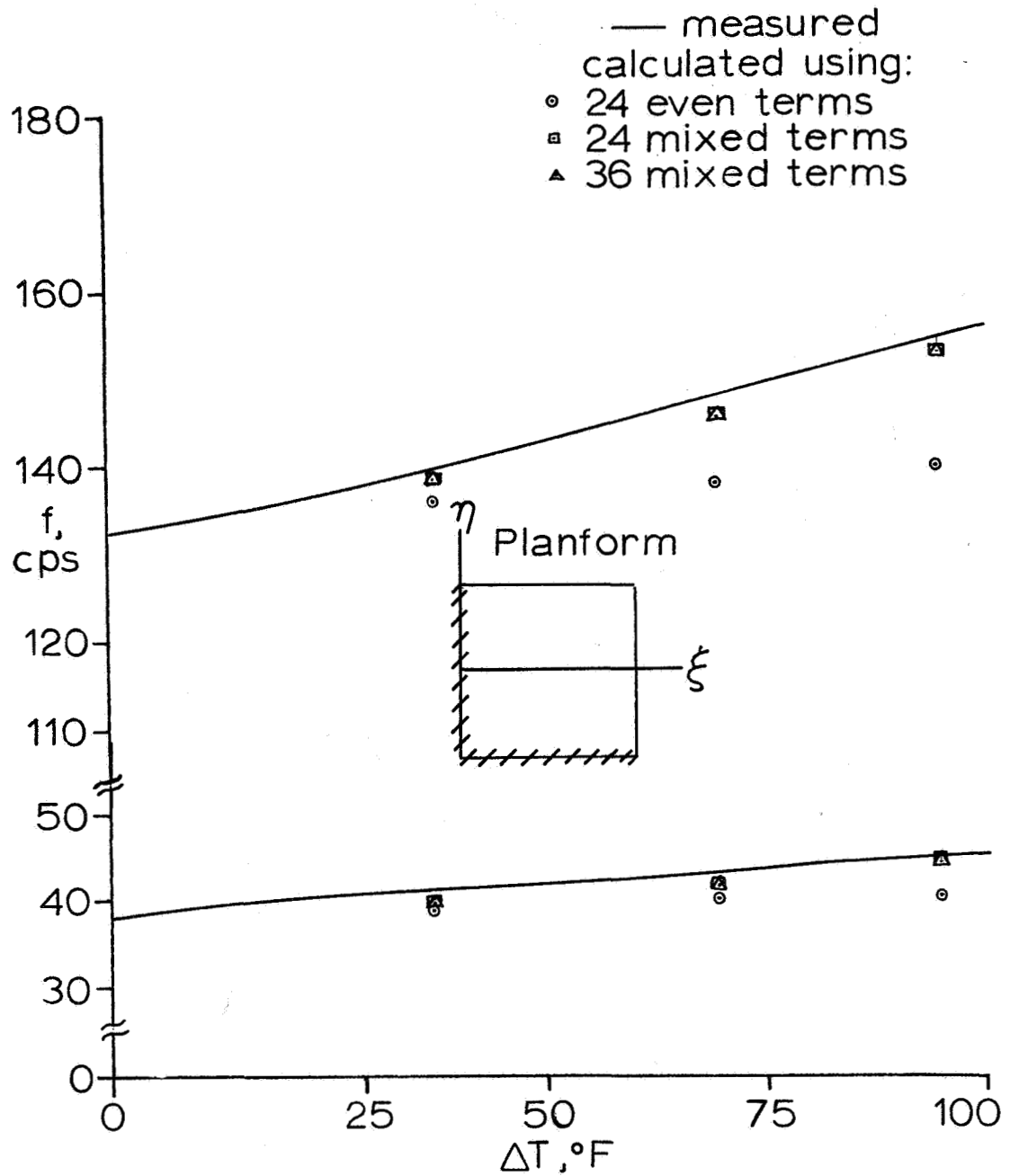


Fig. 18

## Appendix A

### Matrix Elements and Parameters

$$\begin{aligned}
 B_{ij,kl} = & \iint \left( \frac{h}{h_r} \right)^3 \{ (\alpha_{ij})_{\xi\xi} (\alpha_{kl})_{\xi\xi} + \frac{E_{22}}{E_{11}} \left( \frac{a}{b} \right)^4 (\alpha_{ij})_{\eta\eta} (\alpha_{kl})_{\eta\eta} \\
 & + \left( \frac{a}{b} \right)^2 \frac{E_{12}}{E_{11}} [ (\alpha_{ij})_{\xi\xi} (\alpha_{kl})_{\eta\eta} + (\alpha_{ij})_{\eta\eta} (\alpha_{kl})_{\xi\xi} ] \\
 & + 4 \left( \frac{a}{b} \right)^2 \frac{G_{12}}{E_{11}} (\alpha_{ij})_{\xi\eta} (\alpha_{kl})_{\xi\eta} \} d\xi d\eta
 \end{aligned}$$

$$\begin{aligned}
 M_{ij,kl} = & \iint \{ F_{\eta\eta} (\alpha_{ij})_{\xi} (\alpha_{kl})_{\xi} + F_{\xi\xi} (\alpha_{ij})_{\eta} (\alpha_{kl})_{\eta} \\
 & - F_{\xi\eta} [ (\alpha_{ij})_{\xi} (\alpha_{kl})_{\eta} + (\alpha_{ij})_{\eta} (\alpha_{kl})_{\xi} ] \} d\xi d\eta
 \end{aligned}$$

$$T_{ij,kl} = \iint \frac{h}{h_r} (\alpha_{ij}) (\alpha_{kl}) d\xi d\eta$$

$$\begin{aligned}
 A_{pq,rs} = & \iint \frac{h_r}{h} \{ \left( \frac{a}{b} \right)^4 (\gamma_{pq})_{\eta\eta} (\gamma_{rs})_{\eta\eta} + \frac{a_{22}}{a_{11}} (\gamma_{pq})_{\xi\xi} (\gamma_{rs})_{\xi\xi} \\
 & + \frac{a_{12}}{a_{11}} \left( \frac{a}{b} \right)^2 [ (\gamma_{pq})_{\xi\xi} (\gamma_{rs})_{\eta\eta} + (\gamma_{pq})_{\eta\eta} (\gamma_{rs})_{\xi\xi} ] \\
 & + \frac{b_{12}}{a_{11}} \left( \frac{a}{b} \right)^2 (\gamma_{pq})_{\xi\eta} (\gamma_{rs})_{\xi\eta} \} d\xi d\eta
 \end{aligned}$$

$$\Gamma_{rs} = \iint \left[ \left( \frac{a}{b} \right)^2 (\gamma_{rs})_{\eta\eta} + \frac{\alpha_2}{\alpha_1} (\gamma_{rs})_{\xi\xi} \right] T(\xi, \eta) d\xi d\eta$$

$$\lambda^2 = \omega^2 \, 12 \rho a^4 / E_{11} h_r^2$$

$$k_1 = \frac{12}{E_{11}} \left(\frac{a}{b}\right)^2 \left(\frac{a}{h_r}\right)^2$$

$$k_2 = (\alpha_1 \quad \Delta T/a_{11})$$

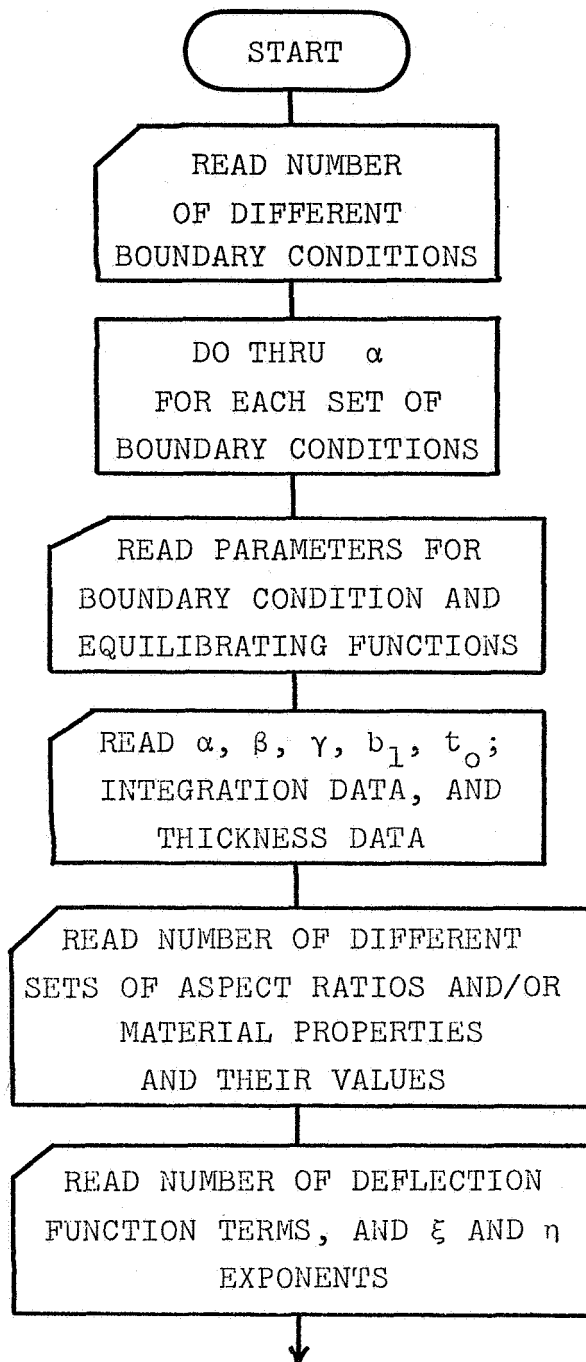
$$\{\hat{C}\} = \frac{1}{a^2 h_r} \{C\}$$

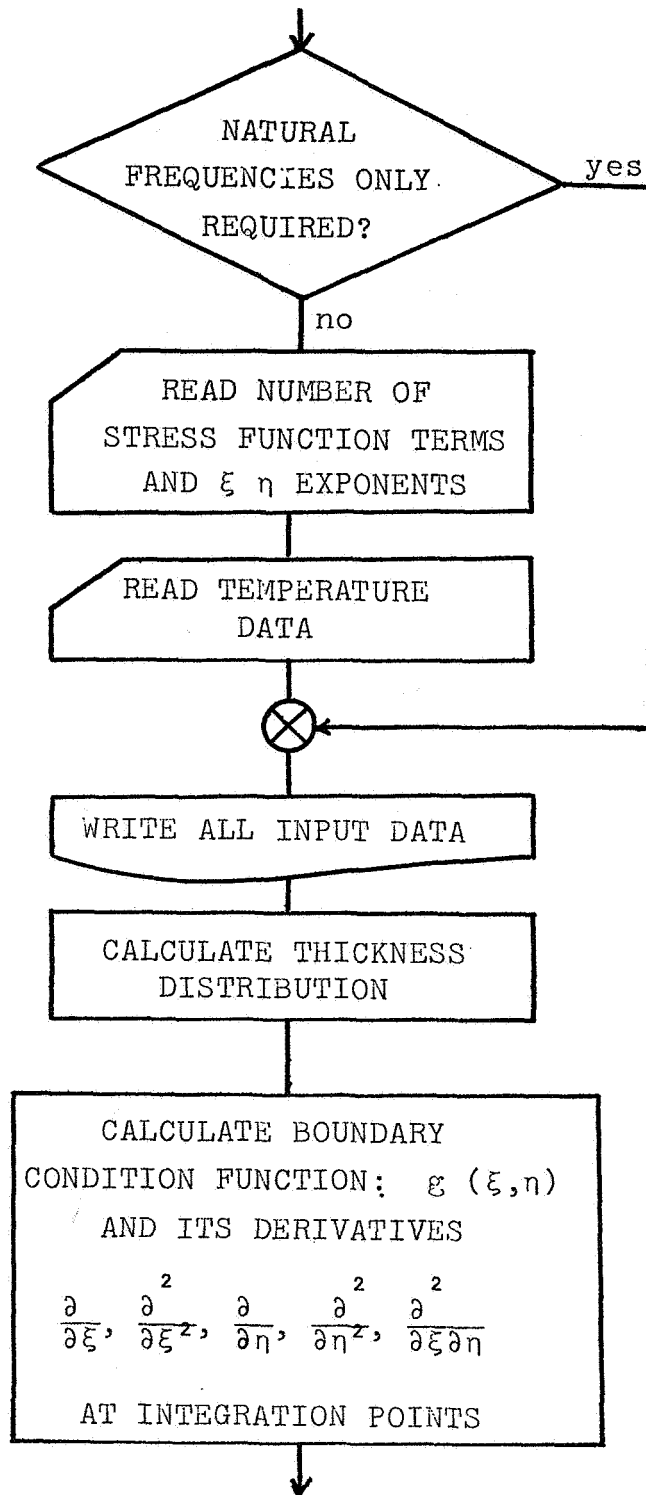
$$\alpha_{ij} = g(\xi,\eta) \, \xi^i \eta^j$$

$$\gamma_{pq} = f(\xi,\eta) \xi^p \eta^q$$

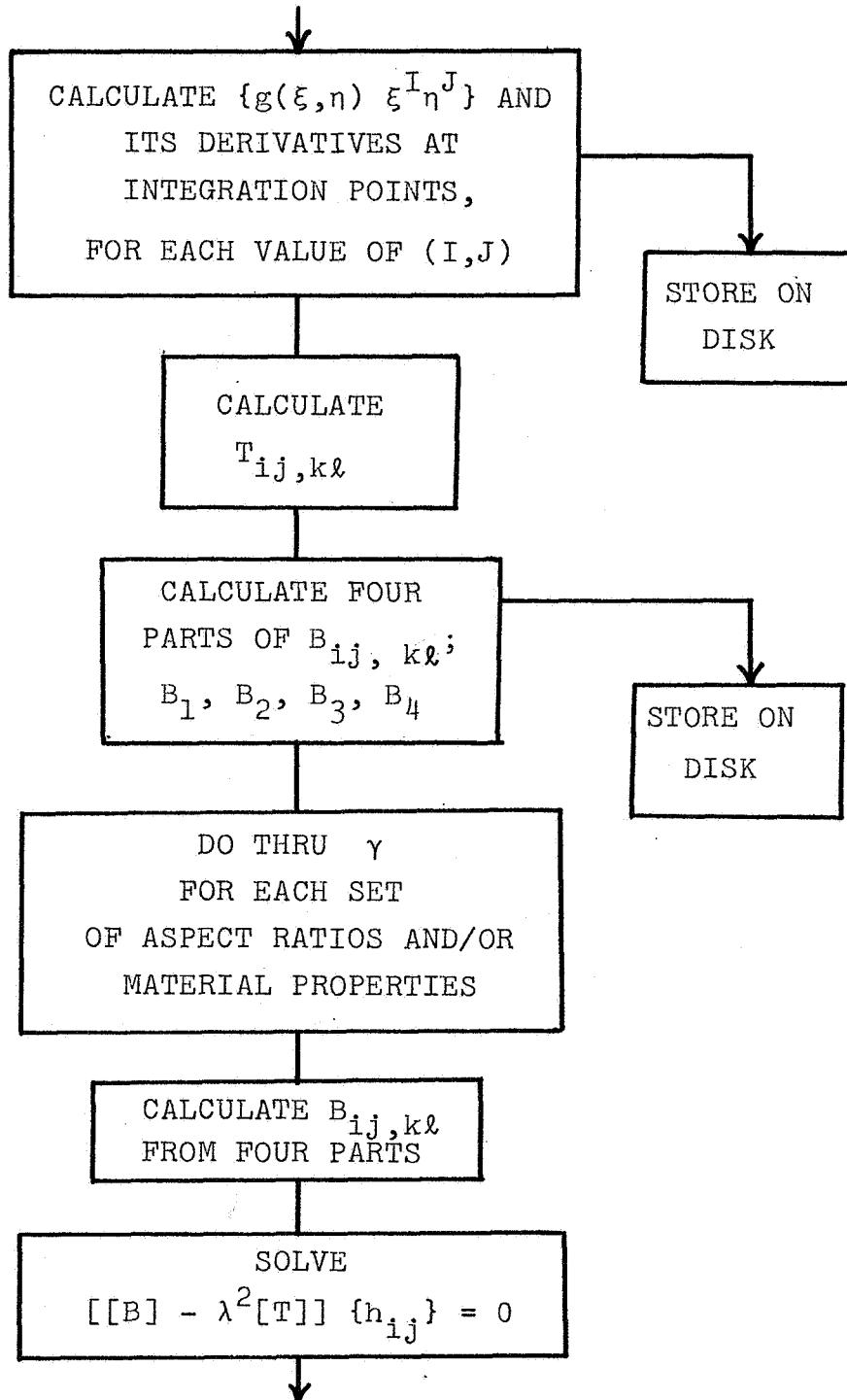
## Appendix B

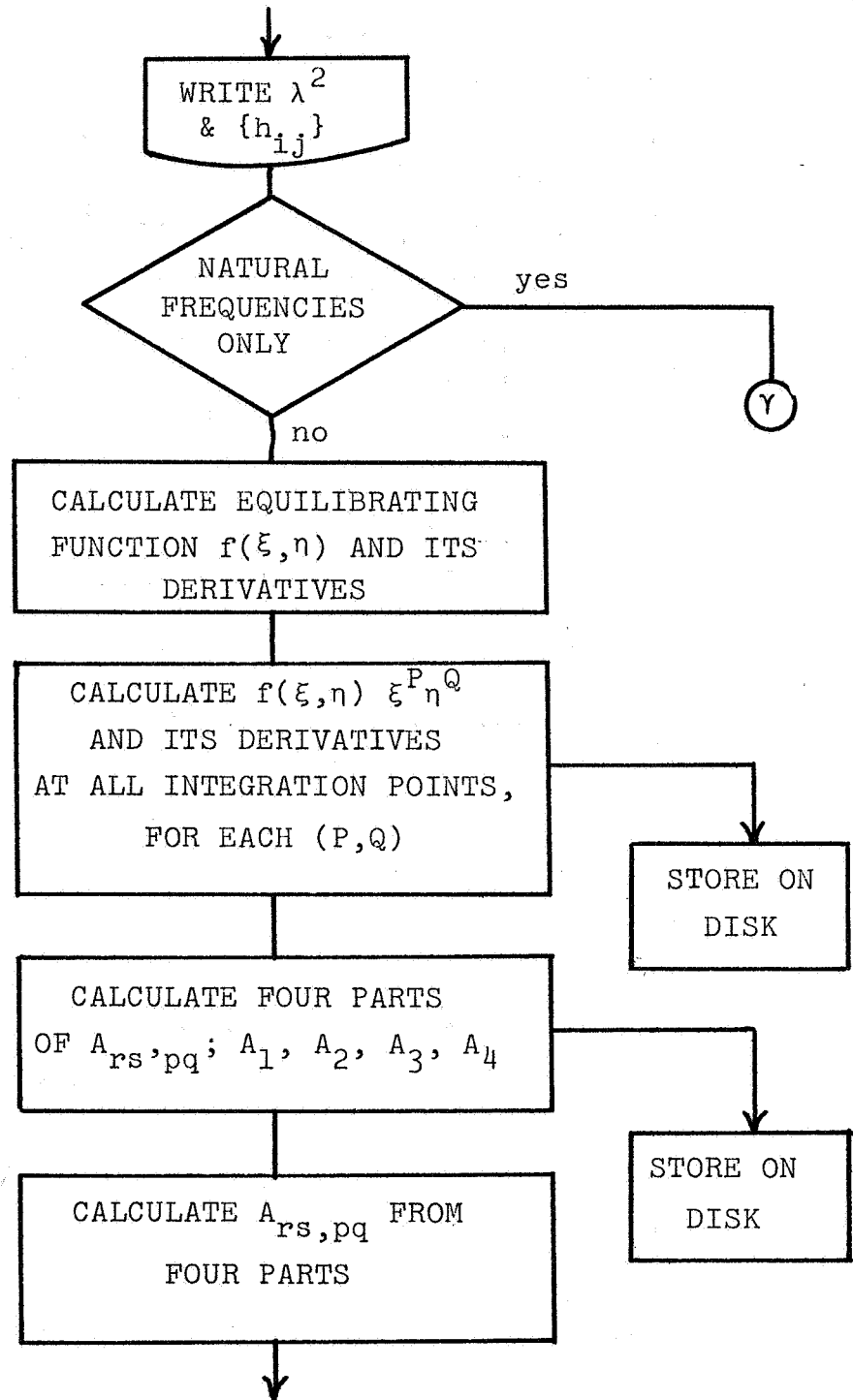
### Logic Flow Diagram

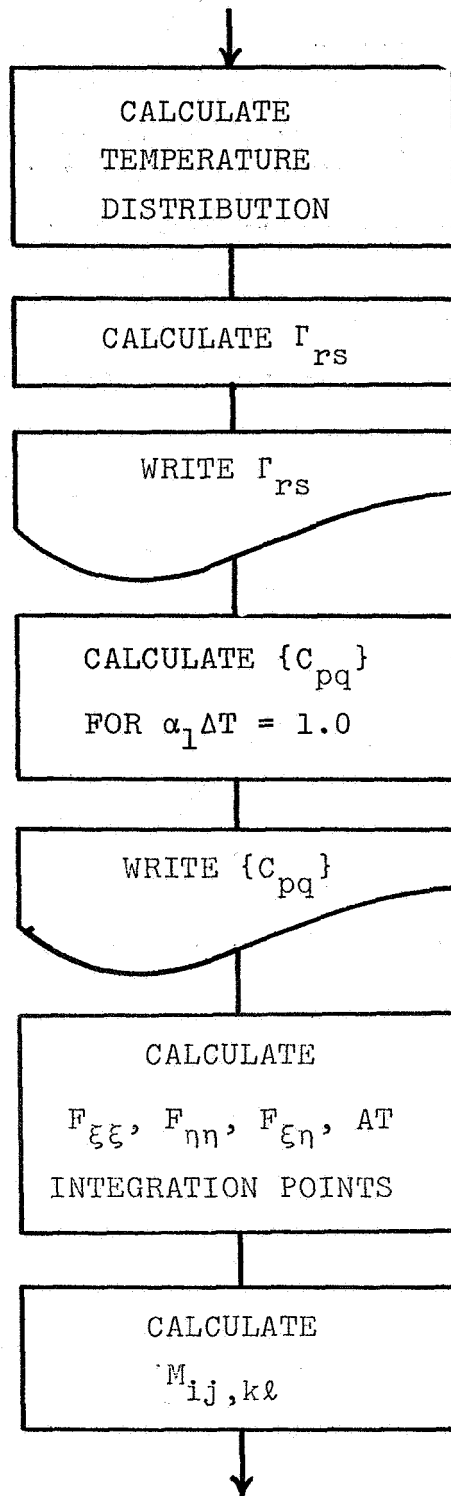


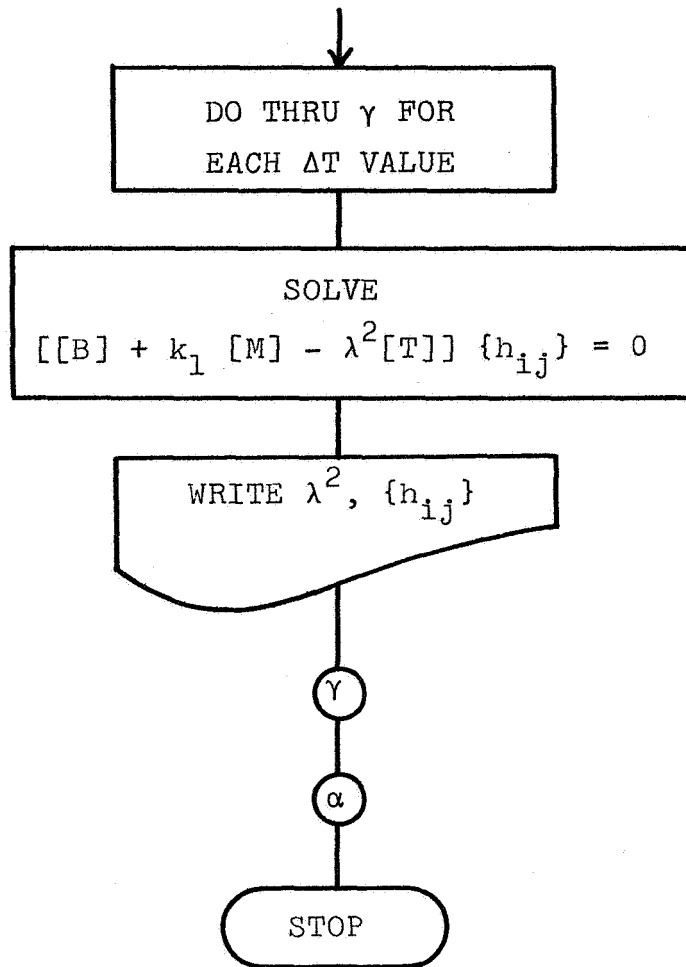












## Appendix C

### Program Listing



```

*13,' QUADRATURE PCINTS')
100 FORMAT(1H /30X,'ASPECT RATIO =' ,E17.8/(45X,E16.8))
130 FCRMAT(1H /16X,'I= ' ,16I4/(19X,16I4))
140 FORMAT(1H /15X,'J= ' ,16I4/(19X,16I4))
150 FCRMAT(1H /16X,'P= ' ,16I4/(15X,16I4))
160 FCRMAT(1H /15X,'Q= ' ,16I4/(19X,16I4))
170 FORMAT(1H /37X,'THICKNESS DISTRIBUTION' ,/,23X,'COEFFICIENTS',5X,
  *X EXPONENTS',5X,'Y EXPONENTS' ,/,47X,'SURFACE 1' ,/, (19X,E16.8,7X,
  *13,11X,I3))
180 FORMAT(1H /46X,'SURFACE 2' ,/, (19X,E16.8,7X,I3,11X,I3))
190 FCRMAT(10A8)
200 FORMAT(1H /37X,'TEMPERATURE DISTRIBUTION' ,/,20X,'COEFFICIENTS',5X,
  *X EXPONENTS',5X,'Y EXPONENTS' ,/, (16X,E16.8,7X,I3,11X,I3))
210 FCRMAT(1H /20X,10A8)
220 FORMAT(1H /30X,'DELTA-T =' ,E17.8/(40X,E16.8))
230 FCRMAT(1H /25X,'INITIAL DEFLECTION COEFFICIENTS'/( 20X,5E16.8))
240 FCRMAT(1H /25X,'TRANSVERSE LOADING COEFFICIENTS'/( 20X,5E16.8))
250 FORMAT(1H,42X,'ARRAY T')
260 FCRMAT(1H /5X,'ROW',I3/(10X,5E16.8))
270 FORMAT(1H,35X,'ASPECT RATIO =' ,E16.8/,42X,'ARRAY B')
280 FORMAT(1H,10X,'FUNDAMENTAL VIBRATION EIGENVALUES SQUARED - ASPECT
  * RATIC=' ,E16.8
  /(10X,5E18.8))
290 FCRMAT(1H /10X,'VIBRATION EIGENVECTORS')
300 FORMAT(1H,35X,'ASPECT RATIO =' ,E16.8/42X,'ARRAY A')
310 FCRMAT(1H,37X,'THERMAL LOADING'/(10X,5E16.8))
320 FORMAT(//30X,'STRESS FUNCTION CCEFFICIENTS'/(10X,5E16.8))
330 FORMAT(1H,42X,'ARRAY M')
340 FORMAT(1H,30X,'LINEAR VIBRATION EIGENVALUES SQUARED a DELTA-T= ',
  *E16.8
  /(10X,5E18.8))
350 FCRMAT(1H /25X,'VIBRATION EIGENVECTORS')
360 FORMAT(1H,56X,A1,' ',A1/55X,'-----'/55X,'| |'/30X,'BOUNDARY
  * CONDITIONS: ',A1,' ',A1,'| |',A1,' ',A1/55X,'| |'/55X,
  * '-----'/57X,A1,' ',A1)
400 FORMAT(1H,12X,'STRESS VARIATION AT X=' ,E13.5/12X,'Y',15X,' NX
  * ',13X,' NY ',13X,' NXY ',13X,'THICK',13X,'TEMP')
410 FORMAT(1X,6E19.8)

```

```

480 FORMAT(1H1,'X')
490 FORMAT(1H1,20X,'TEMPERATURE DISTRIBUTION NO.',I3,' TREF= ',
      *E13.5//)
500 FORMAT(10X,5E15.5)
510 FCRMAT(1H1,//////////,40X,'TEMPERATURE DISTRIBUTION NO.',
      *I3,' OF ',I3)
560 FORMAT(1H,11X,'E11',14X,'E12',14X,'E22',14X,'G12',12X,'ALPHA1',
      *11X,'ALPHA2'/(4X,6E17.7))
      READ(5,20)NBC
      CC 1000 NCCND=1,NBC
      REWIND 2
      REWIND 3
      REWIND 4
      READ(5,30)(BC(I),I=1,4)
      BC(I) GIVES THE DISPLACEMENT BOUNDARY CONDITIONS
      BC(I)= P - SIMPLY SUPPORTED
              C - CLAMPED
              F - FREE
      READ(5,30) (RES(I),I=1,4)
      RES(I) GIVES THE STRESS BOUNDARY CONDITIONS
      RES(I)= C - CLAMPED
              F - FREE
      EDGES ARE NUMBERED CLOCKWISE STARTING WITH THE EDGE CONTAINING
      THE ORIGIN
      READ(5,10) ALPHA,BETA,GAMMA,ACH,B1,TO
      PLATE GEOMETRIC PARAMETERS
      READ(5,20)NX2
      NX2= 1/2 THE NUMBER OF QUADRATURE POINTS
      READ(5,40)(HKER(I),I=1,NX2)
      HKER(I)= QUADRATURE COEFFICIENTS IN ASCENDING ORDER
      READ(5,40)(ZKER(I),I=1,NX2)
      ZKER(I)= QUADRATURE POINTS IN DESCENDING ORDER
      READ(5,20)NTHIC1
      NTHIC1= NUMBER OF THICKNESS TERMS ON SURFACE 1
      READ(5,10)(TCC1(I),I=1,NTHIC1)
      TCC1= THICKNESS FUNCTION COEFFICIENTS ON SURFACE 1

```

C C C C C C C C C C C C C C C C



```

C      READ(5,20)(NTX1(I),I=1,NTHIC1)
C      NTX1= X-EXPONENTS OF THICKNESS FUNCTION ON SURFACE 1
C      READ(5,20)(NTY1(I),I=1,NTHIC1)
C      NTY1= Y-EXPONENTS OF THICKNESS FUNCTION ON SURFACE 1
C      READ(5,20)NTHIC2
C      READ(5,10)(TCO2(I),I=1,NTHIC2)
C      READ(5,20)(NTX2(I),I=1,NTHIC2)
C      READ(5,20)(NTY2(I),I=1,NTHIC2)
C      NTHIC2, TCO2, NTX2, NTY2 ARE SAME AS ABOVE BUT ON SURFACE 2
C      REAC(5,20)NAR
C      NAR = NUMBER OF DIFFERENT SETS OF ASPECT RATIO AND MATERIAL
C      PROPERTIES.
C      DO 2 J=1,NAR
C      READ(5,10) E11(J),E12(J),E22(J),G12(J),AL1(J),AL2(J)
C      MATERIAL PROPERTIES
C      FOR AN ISOTROPIC MATERIAL ;
C      E11=E22=E/(1-NU**2)
C      E12=NU*E/(1-NU**2)
C      G12=E/2*(1+NU)
C      AL1=AL2=AL - (THERMAL EXPANSION COEFF.)
C
C      2 READ(5,10) AR(J)
C      ASPECT RATIO
C      READ(5,20)NDEFL
C      NUMBER OF DEFLECTION FUNCTION TERMS
C      READ(5,20)(IC(I),I=1,NCEFL)
C      X-EXPONENTS OF DEFLECTION FUNCTION
C      REAC(5,20)(JC(I),I=1,NDEFL)
C      Y-EXPONENTS OF DEFLECTION FUNCTION
C      READ(5,50)LAMDAO
C      LAMDAO IS A LOGICAL VARIABLE: = T - END OF INPUT AND ONLY
C      FUNDAMENTAL FREQUENCIES ARE
C      CALCULATED.
C      = F - INPUT CONTINUES.
C
C      IF(LAMDAO) GO TO 3

```

```

C      NUMBER OF STRESS FUNCTION TERMS
      READ(5,20)NSTRES
      READ(5,20)(IP(I),I=1,NSTRES)
C      X-EXPCNENTS OF STRESS FUNCTION
      READ(5,20)(IQ(I),I=1,NSTRES)
C      Y-EXPCNENTS OF STRESS FUNCTION
      READ(5,50)EXPT
C      EXPT IS A LOGICAL VARIABLE: = T - INPUT UP TO 5 EXPERIMENTAL
C      TEMPERATURE DISTRIBUTIONS.
C      F - INPUT 1 ANALYTICAL DISTRIBUTION.

      IF(EXPT) GO TO 4
      NTEMP=1
      READ(5,60)NTEM,TREF(1)
C      NTEM= NUMBER OF TERMS IN TEMPERATURE POLYNOMIAL,
C      TREF= REFERENCE TEMPERATURE
      READ(5,10)(TEM(I),I=1,NTEM)
C      COEFFICIENTS OF TEMPERATURE POLYNOMIAL
      READ(5,20)(NTEMX(I),I=1,NTEM)
C      X-EXPCNENTS OF TEMPERATURE POLYNOMIAL
      READ(5,20)(NTEMY(I),I=1,NTEM)
C      Y-EXPCNENTS OF TEMPERATURE POLYNOMIAL
      READ(5,20)NDT(1)
C      NUMBER OF DELTA-T'S
      NT=NCT(1)
      DO 41 J=1,NT
41  READ(5,10) DT(J,1)
C      DT= VALUES OF DELTA-T
      GO TO 3
4  READ(5,60) NTX,DTX,DTY,XT1,YT1
C      INPUTS FOR SUBROUTINE 'INTP'
      READ(5,20)(KC(I),I=1,NTX)
C      INPUTS FOR SUBROUTINE 'INTP'
      READ(5,20)(LC(I),I=1,NTX)
C      INPUTS FOR SUBROUTINE 'INTP'
      NPTS=0
      DO 6 I=1,NTX

```

```

6  NPTS=NPTS+LC(I)-KC(I)+1
  READ(5,150)(TITLE(I),I=1,10)
  TITLE= SOME DESCRIPTIVE INFORMATION ABOUT THE TEMPERATURES INPUT.
  READ(5,20)NTEMP
  DO 7 I=1,NTEMP
    READ(5,10)TREF(I)
    REFERENCE TEMP. FOR THE I'TH DISTRIBUTION
  READ(5,10)(TEMP(J,I),J=1,NPTS)
  VALUE OF THE I'TH TEMP. DIST. AT EACH OF THE GRID POINTS.
  READ(5,20)NDT(I)
  NUMBER OF DELTA-T'S TO BE CONSIDERED FOR I'TH TEMP. DIST.
  NT=NDT(I)
  DO 7 J=1,NT
    READ(5,10) DT(J,I)
    VALUES OF DELTA-T FOR I'TH DISTRIBUTION
  3  CONTINUE
  NX=2*NX2
  NX4=2*NX*NX
  WRITE(6,360) BC(2),RES(2),BC(1),RES(1),BC(3),RES(3),BC(4),RES(4)
  WRITE(6,80)
  WRITE(6,100) (AR(I),I=1,NAR)
  WRITE(6,560) (E11(I),E12(I),E22(I),G12(I),AL1(I),AL2(I),I=1,NAR)
  WRITE(6,130) (IC(I),I=1,NDEFL)
  WRITE(6,140) (JC(I),I=1,NDEFL)
  WRITE(6,170) (TC01(I),NTX1(I),NTY1(I),I=1,NTHIC1)
  WRITE(6,180) (TCC2(I),NTX2(I),NTY2(I),I=1,NTHIC2)
  IF(LAMDAO) GO TO 8
  WRITE(6,150) (IP(I),I=1,NSTRES)
  WRITE(6,160) (IQ(I),I=1,NSTRES)
  IF(EXPT) GO TO 15
  WRITE(6,200) (TEM(I),NTEMX(I),NTEMY(I),I=1,NTEM)
  GO TO 16
15  WRITE(6,210)(TITLE(I),I=1,10)
16  CCNTINUE
  IF(EXPT) GO TO 12
  DC 32 J=1,NTEMP

```

```

      NT=NDT(J)
32  WRITE(6,220) (DT(I,J),I=1,NT)
12  CONTINUE
14  CONTINUE
      WRITE(6,480)
      IF (.NOT.EXPT) GO TO 8
      DO 11 J=1,NTEMP
      WRITE(6,490) J,TREF(J)
      NT=NDT(J)
      WRITE(6,220) (DT(I,J),I=1,NT)
11  WRITE(6,500) (TEMP(I,J),I=1,NPTS)
      8 CONTINUE

C
C
C      TRANSFORM QUADRATURE COEFFICIENTS AND POINTS TO OUR COORDINATE SYSTEM

      DO 17 I=1,NX2
      HKER(I)=HKER(I)/TWO
      J=NX-I+1
      HKER(J)=HKER(I)
      ZKER(I)=ONE-((ZKER(I)+CNE)/TWO)
17  ZKER(J)=ONE-ZKER(I)

C
C
C      CALCULATE INTEGRATION POINTS, AND THICKNESS @ POINTS

      I1=0
      DO 18 I=1,NX
      XB=BETA*ZKER(I)
      AA=ONE-ZKER(I)
      QP=ONE+ALPHA*ZKER(I)
      CC=B1-GAMMA*ZKER(I)
      DC 18 NSEC=1,2
      DO 18 J=1,NX
      J1=NX+1-J
      I1=I1+1
      ETA(I1)=(QP-XB)*ZKER(J1)+XB
      IF(NSEC.EQ.2) ETA(I1)=-((XB+CD)*ZKER(J)+XB)

```



```

MAT2(I1) =FX*YO+Y0+F*IC(IJ)*X1*YO
MAT3(I1) =(FXX*YO+TWO*IC(IJ)*X1*FX+F*IC(IJ)*(IC(IJ)-1)*X2)*YO
MAT4(I1) =FY*YO+Y0+F*JC(IJ)*XO*Y1
MAT5(I1) =(FYY*Y0+TWC*JC(IJ)*Y1*FY+F*JC(IJ)*(JC(IJ)-1)*Y2)*XO
MAT6(I1) =FXY*XC*YO+FX*JC(IJ)*XC*Y1+FY*IC(IJ)*X1*Y0+F*JC(IJ)*
      *IC(IJ)*X1*Y1
19 CONTINUE
29 WRITE(2) (MAT1(I),MAT2(I),MAT3(I),MAT4(I),MAT5(I),MAT6(I),I=1,NX4)
C
C
C
C
      FT02F001
      =WIJ
END FILE 2
REWIND 2
      READ FROM FT02F001
      CALCULATE FOUR PARTS OF B MATRIX, AND T MATRIX
C
C
C
DC 21 IJ=1,NDEFL
READ(2)(M1(I),M2(I),M3(I),M4(I),M5(I),M6(I),I=1,NX4)
REWIND 2
DO 21 KL=1,IJ
READ(2)(MAT1(I),MAT2(I),MAT3(I),MAT4(I),MAT5(I),MAT6(I),I=1,NX4)
I3=(KL-1)*NDEFL+IJ
I4=(IJ-1)*NDEFL+KL
I1=0
TAAD=ZERO
BAAD=ZERO
BAAC2=ZERO
BAAC3=ZERO
BAAD4=ZERO
DO 22 I=1,NX
TAD1=ZERO
BAD1=ZERO
PAC2=ZERO
RAD3=ZERO
BAD4=ZERO
X=ZKER(I)
DO 22 NSEC=1,2
TAD=ZERO

```

```

BADC1=ZERC
BADD2=ZERO
RACC3=ZERO
BADD4=ZERC
DY=ONE+(ALPHA-BETA)*X
IF(NSEC EC 2) DY=PI+(BETA-GAMMA)*X
DO 24 J=1,NX
  I1=I1+1
  Y=ETA(I1)
  TGRAND=T(I1)*MAT1(I1)*M1(I1)
  BGRAN1=(M3(I1)*MAT3(I1))*T(I1)**3
  RGRAN2=(M5(I1)*MAT5(I1))*T(I1)**3
  BGRAN3=(M3(I1)*MAT5(I1)+M5(I1)*MAT3(I1))*T(I1)**3
  BGRAN4=M6(I1)*MAT6(I1)*T(I1)**3
  TAD=TAD+TGRAND*HKER(J)*DY
  BADD1=BADD1+BGRAN1*HKER(J)*DY
  BACC2=BACC2+RGRAN2*HKER(J)*DY
  BADD3=BADD3+BGRAN3*HKER(J)*DY
  BADD4=BADD4+BGRAN4*HKER(J)*DY
24  TAD1=TAD+TAD1
  BAD1=BAD1+BADD1
  BAC2=BAC2+RACC2
  BAD3=BAC3+RACC3
23  BAD4=BAD4+BADD4
  TAAC=TAAC+TAD1*HKER(I)
  BAAD=BAAC+BAC1*HKER(I)
  RAAD2=BAAD2+BAD2*HKER(I)
  RAAC3=BAAD3+BAD3*HKER(I)
22  BAAD4=BAAC4+BAD4*HKER(I)
  TMA1(T13)=TAAD
  TMA1(T14)=TAAC
  ARRAY1(I3)=BAAD
  ARRAY1(I4)=BAAD
  ARRAY2(I4)=BAAC2
  ARRAY2(I3)=BAAD2
  ARRAY3(I3)=BAAD3

```

```

ARRAY3(I4)=BAAD3
ARRAY4(I3)=BAAD4
21 ARRAY4(I4)=BAAD4
NDF2=NDEF1**2
IF(NAR.EQ.1) GO TO 34
WRITE(4)(ARRAY1(I),ARRAY2(I),ARRAY3(I),ARRAY4(I),I=1,NDF2)
34 CONTINUE
CC 999 NARS=1,NAR
ACB=AR(NARS)*(ONE+BI-(GAMMA-ALPHA)/TWO)
AOB2=AOB**2
AOB4=AOB**4

```

FT04F001

C  
C  
C

```

CALCULATE B MATRIX FROM THE FOUR PARTS PREVIOUSLY CALCULATED

```

```

IF(NARS.EQ.1) GC TO 33
REWIND 4
READ(4)(ARRAY1(I),ARRAY2(I),ARRAY3(I),ARRAY4(I),I=1,NDF2)
33 CONTINUE
Z2=AOB4*E22(NARS)/E11(NARS)
Z3=ACB2*E12(NARS)/E11(NARS)
Z4=ACB2*4.000*G12(NARS)/E11(NARS)
DO 26 I=1,NDF2
26 BMAT(I)=ARRAY1(I)+Z2*ARRAY2(I)+Z3*ARRAY3(I)+Z4*ARRAY4(I)

```

C  
C  
C

```

CALCULATE THE FUNDAMENTAL FREQUENCIES AND MODE SHAPES

```

```

MODE=C
DO 35 I=1,NDF2
ARRAY5(I)=BMAT(I)
35 ARRAY2(I)=TMAT(I)
CALL SING(ARRAY5,ARRAY2,NDEF1,ZERO,M1,ARRAY6,SIN)
IF(SIN) GC TO 145
CALL DNROOT(NDEF1,ARRAY2,ARRAY5,M1,ARRAY6,MODE)
DO 148 I=1,NDEF1
148 M1(I)=ONE/M1(I)
145 WRITE(6,280) AR(NARS),(M1(I),I=1,NDEF1)

```





FT03F001

```

M3(I1)=FPQXY
53 CCNTINUE
43 WRITE(3) (M1(I),M2(I),M3(I),I=1,NX4)
END FILE 3

C      NOW CALCULATE THE FOUR PARTS OF THE A MATRIX
C
C
REWIND 3
DO 54 NM=1,NSTRES
  READ(3) (MAT1(I),MAT2(I),MAT3(I),I=1,NX4)
  REWIND 3
  DO 54 N=1,NM
    READ(3) (M1(I),M2(I),M3(I),I=1,NX4)
    I3=(N-1)*NSTRES+NM
    I4=(NM-1)*NSTRES+N
    AAC1=ZERO
    AAC2=ZERO
    AAC3=ZERO
    AAC4=ZERO
    I1=0
    DO 55 I=1,NX
      EAC1=ZERO
      EAC2=ZERO
      EAC3=ZERO
      EAC4=ZERO
    DO 56 NSEC=1,2
      FAD1=ZERO
      FAD2=ZERO
      FAD3=ZERO
      FAD4=ZERO
      CY=CNE+(ALPHA-BETA)*ZKER(I)
      IF(NSEC.EQ.2)DY=B1+(BETA-GAMMA)*ZKER(I)
      DO 57 J=1,NX
        I1=I1+1
        EGRAN1=MAT1(I1)*M1(I1)/T(I1)
        EGRAN2=MAT2(I1)*M2(I1)/T(I1)

```

```

EGRAN3=(MAT1(I1)*M2(I1)+M1(I1)*MAT2(I1))/T(I1)
EGRAN4=(MAT3(I1)*M3(I1))/T(I1)
FAD1=FAC1+EGRAN1*HKER(J)*DY
FAD2=FAD2+EGRAN2*HKER(J)*DY
FAC3=FAC3+EGRAN3*HKER(J)*DY
57 FAD4=FAD4+EGRAN4*HKER(J)*DY
EAD1=EAD1+FAD1
EAC2=EAC2+FAC2
EAD3=EAD3+FAD3
56 EAD4=EAD4+FAD4
AAC1=AAD1+EAD1*HKER(I)
AAD2=AAD2+EAD2*HKER(I)
AAD3=AAD3+EAD3*HKER(I)
55 AAD4=AAD4+EAD4*HKER(I)
ARRAY1(I3)=AAD1
ARRAY1(I4)=AAC1
ARRAY2(I4)=AAC2
ARRAY2(I3)=AAD2
ARRAY3(I4)=-AAD3
ARRAY3(I3)=-AAD3
ARRAY4(I3)=AAD4
ARRAY4(I4)=AAC4
54 CONTINUE

```

FT04F001

```

C
C   STORE THESE MATRICES IN DSRN=4
NSTRS2=NSTRS**2
IF(NAR.EQ.1) GO TO 51
WRITE(4) (ARRAY1(I),ARRAY2(I),ARRAY3(I),ARRAY4(I),I=1,NSTRS2)
END FILE 4
51 CONTINUE
*   IF(NARS.NE.1) READ(4) (ARRAY1(I),ARRAY2(I),ARRAY3(I),ARRAY4(I),
      I=1,NSTRS2)
C
C   CALCULATE A MATRIX FROM ITS FOUR PARTS
Y1=E11(NARS)/E22(NARS)
Y2=AQB4

```

```

Y3=AOB2*E12(NARS)/E22(NARS)
Y4=AOB2*(E11(NARS)*E22(NARS)-E12(NARS)**2)/(G12(NARS)*E22(NARS))
DC 58NM=1,NSTRES
DO 58 N=1,NM
I3=(N-1)*NSTRES+NM
I4=(NM-1)*NSTRES+N
AMAT(I3)=Y1*ARRAY1(I3)+Y2*ARRAY2(I3)+Y3*ARRAY3(I3)+Y4*ARRAY4(I3)
58 AMAT(I4)=AMAT(I3)

```

C  
C  
C

# CALCULATE TEMPERATURE AT QUADRATURE POINTS

```

DO 598 NTS=1,NTEMP
WRITE(6,510) NTS,NTEMP
N2X=NX*2
IF(EXPT) GO TO 62
I1=0
DO 63 I=1,NX
X=ZKER(I)
DO 63 J=1,N2X
I1=I1+1
Y=ETA(I1)
63 TEMPT(I1)=CTEM(NTEM ,TEM,NTEMX,NTEMY,X,Y,ALPHA,B1,GAMMA)
GO TO 64
62 CONTINUE
DC 65 I=1,NPTS
65 MAT1(I)=TEMP(I,NTS)
I1=0
DC 66 I=1,NX
X=ZKER(I)
DC 66 J=1,N2X
I1=I1+1
Y=ETA(I1)
CALL INTP(MAT1,KC,LC,X,Y,DTY,DTX,NTX,YT1,XT1,WS,WXS)
66 TEMPT(I1)=WS
64 CONTINUE

```

C

```

C          CALCULATE GAMRS - THERMAL LOADING
REWIND 3
DC 67NM=1,NSTRES
READ(3) (M1(I),M2(I),M3(I),I=1,NX4)
I1=0
GAD1=ZERO
DO 68 I=1,NX
  GAD2=ZERO
  X=ZKER(I)
  DO 69 NSEC=1,2
    GAD3=ZERO
    DY=ONE+(ALPHA-BETA)*X
    IF (NSEC.EQ. 2) DY=B1+(BETA-GAMMA)*X
    DO 71 J=1,NX
      I1=I1+1
      Y=ETA(I1)
      71 GAD3=GAD3+(AL2(NARS)/ALL(NARS)*M1(I1)+ADB2*M2(I1))*TEMPT(I1)*
        *HKER(J)*DY
      69 GAD2=GAD2+GAD3
      68 GAD1=GAD1+GAD2*HKER(I)
      67 GAMRS(NM)=(E11(NARS)*E22(NARS)-E12(NARS)**2)/E22(NARS)*GAD1
      WRITE(6,310) (GAMRS(I),I=1,NSTRES)

C          CALCULATE CPQ DUE TO TEMPERATURE WITH ALPHA*DT = 1 0
C
C
C          IF (NTS.EQ.1) CALL CMINV(AMAT,NSTRES,D,M1,M2)
C          *AMAT* NOW CONTAINS THE INVERSE OF THE A-MATRIX
C          DO 36 I=1,NSTRES
C            36 GAMRS(I)=-GAMRS(I)
C            CALL DGMPRD(AMAT,GAMRS,CPQ,NSTRES,NSTRES,1)
C            WRITE(6,320) (CPQ(I),I=1,NSTRES)

C          CALCULATE THE THERMAL STRESSES
C
C          REWIND 3
          READ(3) (MAT1(I),MAT2(I),MAT3(I),I=1,NX4)

```

```

DC201 I=1,NX4
M2(I)=MAT1(I)*CPQ(I)
M1(I)=MAT2(I)*CPQ(I)
201 M3(I)=MAT3(I)*CPQ(I)
DO202 IPQ=2,NSTRES
  REAC(3) (MAT1(I),MAT2(I),MAT3(I),I=1,NX4)
DO202 I=1,NX4
  M2(I)=M2(I)+MAT1(I)*CPQ(IPQ)
  M1(I)=M1(I)+MAT2(I)*CPQ(IPQ)
  202 M3(I)=M3(I)+MAT3(I)*CPQ(IPQ)
  A=E22(NARS)/(E11(NARS)*E22(NARS)-E12(NARS)**2)
  I1=C
DO 42 I=1,NX
  WRITE(6,400) ZKER(I)
DO 42 NSEC=1,2
DO 42 J=1,NX
  I1=I1+1
  SXX=M1(I1)*TO*ACB2*A
  SYX=M2(I1)*TO*A
  SXY=-M3(I1)*TO*ACB*A
  42 WRITE(6,410) ETA(I1),SXX,SYX,SXY,T(I1),TEMPT(I1)

      CALCULATE THE 2-D M-MATRIX

REWIND 2
DC 44 IJK=1,NCEFL
READ(2) (MAT2(I),MAT1(I),MAT4(I),MAT3(I),MAT5(I),MAT6(I),I=1,NX4)
REWIND 2
DC 44 IKL=1,IJK
READ(2) (ARRAY1(I),ARRAY2(I),ARRAY3(I),ARRAY4(I),ARRAY5(I),
*ARRAY6(I),I=1,NX4)
DO 45 I=1,NX4
  MAT2(I)=ARRAY2(I)
  45 MAT4(I)=ARRAY4(I)
  I3=(IJK-1)*NDEFL+IKL
  I4=(IKL-1)*NDEFL+IJK

```

C  
C  
C

```

EAD1=ZERC
I1=0
DC 47 I=1,NX
EAD2=ZERC
X=ZKER(I)
DC 48 NSEC=1,2
DY=ONE+(ALPHA-BETA)*X
IF(NSEC EQ 2) DY=B1+(BETA-GAMMA)*X
EAD3=ZERC
DO 49 J=1,NX
I1=I1+1
EGRAND=M1(I1)*MAT2(I1)*MAT1(I1)+M2(I1)*MAT3(I1)*MAT4(I1)-M3(I1)*
*(MAT1(I1)*MAT4(I1)+MAT2(I1)*MAT3(I1))
49 EAD3=EAD3+EGRAND*HKER(J)*DY
48 EAD2=EAD2+EAD3
47 EAC1=EAC1+EAD2*HKER(I)
RMAT(I3)=EAD1
44 RMMAT(I4)=EAD1

NT=NDT(NTS)
DO 997 NDT5=1,NT

SOLVE LINEAR RESPONSE PROBLEM

Q=12 ODO*A0B2*ADT**2*AL1(NARS)*DT(NDTS,NTS)/E11(NARS)
DO 76 I=1,NDF2
ARRAY1(I)=TMAT(I)
76 ARRAY2(I)=BMAT(I)+RMMAT(I)*Q
CALL SING(ARRAY2,ARRAY1,NDEFL,ZERO,M1,ARRAY3,SIN)
IF(SIN) GO TO 146
CALL DNROCT(NDEFL,ARRAY1,ARRAY2,M1,ARRAY3,MODE)
DO 149 I=1,NDEFL
149 M1(I)=ONE/M1(I)
146 FDT=DT(NDTS,NTS)
WRITE(6,340) FDT,(M1(I),I=1,NDEFL)
WRITE(6,350)

```

```
I2=0
DC 78 I=1,NDEFI
I1=I2+1
I2=I2+NDEFI
78 WRITE(6,260) I,(ARRAY3(J),J=I1,I2)
996 CONTINUE
997 CONTINUE
998 CONTINUE
999 CONTINUE
1000 CONTINUE
      STCP
      END
```



```

SUBROUTINE FUNCTN(F,FX,FXX,FY,FYY,FX,Y,X,Y,S, STRESS,NX,A,G,B1) 5
IMPLICIT REAL*8 (A-H,O-Z) 6
DIMENSION F(1),FX(1),FXX(1),FY(1),FYY(1),FXY(1),X(1),Y(1),S(1),Q(1) 10
*,B(4),IEX(4,3),IR(4),IE(4),T(4,3) 15
LOGICAL STRESS 20
C 25
C 30
C 35
C 40
C 45
C 50
C 60
C 65
C 70
C 91

      ( TRUE -STRESS FUNCTION
      ( FALSE -DISPLACEMENT FUNCTION

STRESS IS A LOGICAL VARIABLE=(
      ( TRUE -STRESS FUNCTION
      ( FALSE -DISPLACEMENT FUNCTION

S(1) IS AN ALPHAMERIC VARIABLE:
('C' IF EDGE(1) IS CLAMPED
('P' IF EDGE(1) IS PINNED
('F' IF EDGE(1) IS FREE

DATA P,FF,CC/'P','F','C'/

      CALCULATE EXPONENTS REQUIRED

DO 1 I=1,4
1 IE(I)=0
IF (STRESS) GO TO 2
DC 21 I=1,4
IF(S(1).EQ.CC) IE(I)=2
IF(S(1).EQ.P) IE(I)=1
21 CCNTINUE
GO TO 3
2 CCNTINUE
DC 4 I=1,4
IF(S(1).EQ.FF.OR.S(1).EQ.P) IE(I)=2
4 CCNTINUE
3 CCNTINUE
DO 5 I=1,4
MP=IE(I)
DO 5 J=1,3
IEX(I,J)=MP-(J-1)

```



```

*      -2.000*IEX(2,1)*IEX(4,1)*T(1,1)*T(2,2)*T(3,1)*T(4,2)      FUN 270
*
*      FXY(11)=(IEX(4,1))*((IEX(1,1)*T(1,2)*T(2,1)*T(3,1)*T(4,2))+(A*IEX(2FUN 275
*      *1)*T(1,1)*T(2,2)*T(3,1)*T(4,2))-(IEX(3,1)*T(1,1)*T(2,1)*T(3,2)*T(FUN 280
*      *4,2))-(G*IEX(4,2)*T(1,1)*T(2,1)*T(4,3))-(IEX(2,1)*T(3,1)*T(1FUN 285
*      *1)*T(1,2)*T(2,2)*T(3,1)*T(4,1))+(A*IEX(2,2)*T(1,1)*T(2,3)*T(3,1)*FUN 290
*      *T(4,1))-(IEX(3,1)*T(1,1)*T(2,2)*T(3,2)*T(4,1))-(G*IEX(4,1)*T(1,1)*FUN 295
*      *T(2,2)*T(3,1)*T(4,2)))      FUN 300
*
*      FXX(11)=IEX(1,1)*IEX(1,2)*T(1,3)*T(2,1)*T(3,1)*T(4,1)      FUN 305
*      +A**2*IEX(2,1)*IEX(2,2)*T(1,1)*T(2,3)*T(3,1)*T(4,1)      FUN 310
*      +IEX(3,1)*IEX(3,2)*T(1,1)*T(2,1)*T(3,3)*T(4,1)      FUN 315
*      +G**2*IEX(4,1)*IEX(4,2)*T(1,1)*T(2,1)*T(3,1)*T(4,3)      FUN 320
*      +2.000*A*IEX(1,1)*IEX(2,1)*T(1,2)*T(2,2)*T(3,1)*T(4,1)      FUN 325
*      -2.000*IEX(1,1)*IEX(3,1)*T(1,2)*T(2,1)*T(3,2)*T(4,1)      FUN 330
*      -2.000*IEX(1,1)*IEX(4,1)*T(1,2)*T(2,1)*T(3,1)*T(4,2)*G      FUN 340
*      -2.000*A*IEX(2,1)*IEX(3,1)*T(1,1)*T(2,2)*T(3,2)*T(4,1)      FUN 345
*      -2.000*A*G*IEX(2,1)*IEX(4,1)*T(1,1)*T(2,2)*T(3,1)*T(4,2)      FUN 350
*      +2.000*G*IEX(3,1)*IEX(4,1)*T(1,1)*T(2,1)*T(3,2)*T(4,2)      FUN 385
*
*      7 CONTINUE      FUN 390
*      RETURN      FUN 395
*      END

```

```

SUBROUTINE PQ(IP,IQ,FUN,FUNX,FUNXX,FUNY,FUNYY,XI,ETA,M,II,I,
*SXPQ,SYPQ,SXYPQ)
  IMPLICIT REAL*8 (F,X,E,S,Y)
  DIMENSION FUN(1),FUNX(1),FUNXX(1),FUNY(1),FUNYY(1),XI(1),
*ETA(1),IP(1),IQ(1)
  X2=XI(1)**(IP(M)-2)
  X1=X2*XI(1)
  X0=X1*XI(1)
  Y2=ETA(II)**(IQ(M)-2)
  Y1=Y2*ETA(II)
  Y0=Y1*ETA(II)
  SXPQ=(FUNXX(II)*X0+2.0*IP(M)*X1*FUNX(II)+FUN(II)*IP(M)*(IP(M)-1)*X
*2)*Y0
  SYPQ=(FUNYY(II)*Y0+2.0*FUNY(II)*IQ(M)*Y1+FUN(II)*IQ(M)*(IQ(M)-1)*Y
*2)*XC
  SXYPQ=FUNXY(II)*XC*YC+FUNX(II)*IQ(M)*X0*Y1+FUNY(II)*IP(M)*X1*Y0+FU
*N(II)*IP(M)*IQ(M)*X1*Y1
  RETURN
  END

```

```

DOUBLE PRECISION FUNCTION CTHIC(NTHIC1,TCO1,NTX1,NTY1,NTHIC2,TCO2,
*NTX2,NTY2,X,Y,BETA)
IMPLICIT REAL*8 (X,B,Y,T,C)
DIMENSION TCO1(1),NTX1(1),NTY1(1),TCO2(1),NTX2(1),NTY2(1)
X1=BETA*X
Y1=Y-X1
IF(Y LT X1) GO TO 1
T=0
DO2 I=1,NTHIC1
2 T=T+TCO1(I)*(X**NTX1(I))*(Y1**NTY1(I))
GC TC 4
1 T=0 C
DO3 I=1,NTHIC2
3 T=T+TCO2(I)*(X**NTX2(I))*(Y1**NTY2(I))
4 CTHIC=T
RETURN
END

```

```

DOUBLE PRECISION FUNCTION CTM(NTEMP,TEM,NTEMX,NTEMY,X,Y,ALPHA,B1,
*GAMMA)
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION TEM(1),NTEMX(1),NTEMY(1)
  A1=B1+1 ODO
  YP=(- X*((ALPHA+GAMMA)
      )+2. ODO*Y-1 ODO+B1)/A1
  YP=DABS(YP)
  TADD=C ODO
  DC1 K=1,NTEMP
  1 TADD=TADD+TEM(K)*(YP*NTEMY(K))*(X*NTEMX(K))
  CTM=TADD
  RETURN
END

```

```

SUBROUTINE SING(A1,A2,N,ZERO,EVAL,EVECT,SIN)
IMPLICIT REAL*8 (A-H,C-Z)
DIMENSION A1(1),A2(1),EVAL(1),EVECT(1)
LOGICAL SIN
SIN= FALSE
I2=0
DO 1 I=1,N
  I1=I2+1
  I2=I2+N
  DO 2 J=I1,I2
    IF(A1(J).NE. ZERO) GO TO 1
  2 CONTINUE
  GC TC 3
  1 CONTINUE
  GO TO 4
  3 CALL DARCOT(N,A1,A2,EVAL,EVECT,0)
  SIN= TRUE
  4 CONTINUE
  RETURN
END

```

```

SUBROUTINE DNROOT (M,A,B,XL,X,MODE)
DIMENSION A(1),B(1),XL(1),X(1)
DOUBLE PRECISION A,B,XL,X,SUMV
IF (MODE EQ. 1) GC TO 101
K=1
DC 100 J=2,M
L=M*(J-1)
DO 100 I=1,J
L=L+1
K=K+1
100 B(K)=B(L)
C
C      THE MATRIX B IS A REAL SYMMETRIC MATRIX
C
MV=0
CALL DEIGEN (B,X,M,MV)
C
C      FCRM RECIPROCAL OF SQUARE ROOT OF EIGENVALUES      THE RESULTS
C      ARE PREMULTIPLIED BY THE ASSOCIATED EIGENVECTORS
C
L=0
DO 110 J=1,M
L=L+J
110 XL(J)=1 C/DSQRT(DABS(B(L)))
K=0
DC 115 J=1,M
DO 115 I=1,M
K=K+1
115 B(K)=X(K)*XL(J)
101 CONTINUE
C
C      FCRM (B**(-1/2))PRIME * A * (B**(-1/2))
C
DO 120 I=1,M
N2=0
DO 120 J=1,M

```

NR00 380

NR00 610

NR00 620

NR00 630

NR00 640

NR00 650

NR00 660

NR00 670

NR00 680

NR00 690

NR00 700

NR00 710

NR00 730

NR00 740

NR00 750

NR00 760

NR00 770

NR00 780

NR00 790

NR00 810

NR00 820

NR00 830

NR00 840

NR00 850

NR00 860

NR00 870

NR00 880

NR00 890

NR00 900

NR00 910



```

N1=M*(I-1)
L=M*(J-1)+I
X(L)=0.0
DO 120 K=1,M
  N1=N1+1
  N2=N2+1
  120 X(L)=X(L)+B(N1)*A(N2)
  L=L+1
  DO 130 J=1,M
    DO 130 I=1,J
      N1=I-M
      N2=M*(J-1)
      L=L+1
      A(L)=0.0
      DO 130 K=1,M
        N1=N1+M
        N2=N2+1
        130 A(L)=A(L)+X(N1)*B(N2)
      C
      C
      C
      CCMPUTE EIGENVALUES AND EIGENVECTORS OF A
      CALL CEIGEN (A,X,M,MV)
      L=0
      DO 140 I=1,M
        L=L+I
        140 XL(I)=A(L)
      C
      C
      C
      CCMPUTE THE NORMALIZED EIGENVECTORS
      DO 150 I=1,M
        N2=0
        DO 150 J=1,M
          N1=I-M
          L=M*(J-1)+I
          A(L)=0.0
          DO 150 K=1,M

```

```

NR00 920
NR00 930
NR00 940
NR00 950
NR00 960
NR00 970
NR00 980
NR00 990
NR00 1000
NR00 1010
NR00 1020
NR00 1030
NR00 1040
NR00 1050
NR00 1060
NR00 1070
NR00 1080
NR00 1090
NR00 1100
NR00 1110
NR00 1120

NR00 1140
NR00 1150
NR00 1160
NR00 1170
NR00 1180
NR00 1190
NR00 1200
NR00 1210
NR00 1220
NR00 1230
NR00 1240
NR00 1250
NR00 1260
NR00 1270

```

```

N1=N1+M
N2=N2+1
150 A(L)=A(L)+B(N1)*X(N2)
L=L+1
K=0
DC 180 J=1,M
SUMV=C,C
DC 170 I=1,M
L=L+1
170 SUMV=SUMV+A(L)*A(L)
175 SUMV=DSQRT(SUMV)
DC 180 I=1,M
K=K+1
180 X(K)=A(K)/SUMV
RETURN
END

```

```

NR001280
NR001290
NR001300
NR001310
NR001320
NR001330
NR001340
NR001350
NR001360
NR001370

NR001390
NR001400
NR001410
NR001420
NR001430

```

```

C
SUBROUTINE INTP(W,K,L,YO,XO,DX,DY,NY,XI,YI,WANS,WXANS)
C
C      IMPLICIT REAL*8 (A,B,C,D,E,F,G,H,I,P,Q,R,S,T,U,V,W,X,Y,Z)
C
C      DIMENSION II(4),JJ(4),K(1),L(1),WW(4),WWX(4),W(1)
C      W FUNCTIONAL VALUES ON A CLOSED RECTANGULAR NODAL POINT SET
C      NY NUMBER OF STRIPS IN RECTANGULAR NODAL POINT SET
C      K,L LIMITS DEFINING NODAL POINT BOUNDARIES OF THE NY STRIPS
C      DX,DY SPACING OF THE NODAL POINTS
C      XI,YI COORDINATES OF REFERENCE NODAL POINT IN XC,YO SYSTEM
C      WANS FUNCTIONAL VALUE AT (XC,YO)
C      WXANS FUNCTIONAL DERIVATIVE OF WANS IN THE XC DIRECTION
C
C      SHIFT COORDINATES TO NODAL POINT COORDINATE SYSTEM
C
C      YY=YO-YI
C      XX=XC-XI
C
C      DETERMINE IDENTIFICATION NUMBERS OF SURROUNDING NODAL POINTS
C
C      12 KMIS=0
C      XI=XX/DX
C      YI=YY/DY
C      II(1)=XI
C      II(1)=II(1)+1
C      II(3)=II(1)
C      II(2)=II(1)+1
C      II(4)=II(2)
C      JJ(1)=YI
C      JJ(1)=JJ(1)+1
C      JJ(2)=JJ(1)
C      JJ(3)=JJ(2)+1
C      JJ(4)=JJ(3)
C
INTP0015
INTP0020
INTP0030
INTP0025
INTP0035
INTP0040
INTP0045
INTP0050
INTP0055
INTP0060
INTP0065
INTP0070
INTP0075
INTP0080
INTP0085
INTP0090
INTP0095
INTP0100
INTP0105
INTP0110
INTP0115
INTP0120
INTP0125
INTP0130
INTP0135
INTP0140
INTP0145
INTP0150
INTP0155
INTP0160
INTP0165
INTP0170
INTP0175
INTP0180
INTP0185

```

INTP0190  
INTP0195  
INTP0200  
INTP0205  
INTP0210  
INTP0215  
INTP0220  
INTP0225  
INTP0230  
INTP0235  
INTP0240  
INTP0245  
INTP0250  
INTP0255  
INTP0260  
INTP0265  
INTP0270  
INTP0275  
INTP0280  
INTP0285  
INTP0290  
INTP0295  
INTP0300  
INTP0305  
INTP0310  
INTP0315  
INTP0320  
INTP0325  
INTP0330  
INTP0335  
INTP0340  
INTP0345  
INTP0350  
INTP0355  
INTP0360  
INTP0365

```

DO 15 M=1,4
JC=JJ(M)
IF(II(M),LT K(JO))GO TO 100
IF(II(M),GT L(JO))GO TO 100
GO TO 16

C
C
C
C
POINT II(M),JJ(M) IS NOT AN INPUT POINT

100 IF(M EQ 1) KMIS=KMIS+1
IF(M EQ 2) KMIS=KMIS+2
IF(M EQ 3) KMIS=KMIS+4
IF(M EQ 4) KMIS=KMIS+8

C
C
C
C
KMIS DETERMINES WHICH SURROUNDING NODAL POINTS ARE PRESENT

GO TO 15

C
C
C
C
POINT II(M),JJ(M) IS AN INPUT POINT

PROCEED NOW TO FIT TWO DIMENSIONAL PARABOLA
ABOUT THE POINTII(M),JJ(M)

16 JCJ=JJ(M)
IOI=II(M)
JM1=JOJ-1
JP1=JCJ+1
IM1=IOI-1
IP1=IOI+1

C
C
C
C
L1 LOCATION OF W(IOI,JM1)
L2 LOCATION OF W(IM1,JCJ)
L3 LOCATION OF W(IOI,JOJ)

```

INTP0370  
 INTP0375  
 INTP0380  
 INTP0385  
 INTP0390  
 INTP0395  
 INTP0400  
 INTP0405  
 INTP0410  
 INTP0415  
 INTP0420  
 INTP0425  
 INTP0430  
 INTP0435  
 INTP0440  
 INTP0445  
 INTP0450  
 INTP0455  
 INTP0460  
 INTP0465  
 INTP0470  
 INTP0475  
 INTP0480  
 INTP0485  
 INTP0490  
 INTP0495  
 INTP0500  
 INTP0505  
 INTP0510  
 INTP0515  
 INTP0520  
 INTP0525  
 INTP0530  
 INTP0535  
 INTP0540  
 INTP0545

L4 LOCATION OF W(IP1,JOJ)  
 L5 LOCATION OF W(IOI,JPI)

L3=0  
 IF(JM1,LT,1)GO TO 505  
 DO 19 KK=1,JM1  
 L3=L3+L(KK)-K(KK)  
 L3=L3+IOI+JOJ-K(JOJ)  
 L1=L3+K(JCJ)-L(JM1)-1  
 L5=L3+L(JOJ)-K(JPI)+1  
 L2=L3-1  
 L4=L3+1

19  
 505

IF KB1=0 W(IOI,JM1) NOT AN INPUT POINT  
 IF KB2=0 W(IP1,JOJ) NOT AN INPUT POINT  
 IF KB3=0 W(IM1,JOJ) NOT AN INPUT POINT  
 IF KB4=0 W(IOI,JPI) NOT AN INPUT POINT

IF(JM1,LT,1) GO TO 501  
 KB1=1  
 IF(IOI,LT,K(JM1))KB1=0  
 IF(ICI,GT,L(JM1))KB1=0  
 GO TO 502

501 KB1=0  
 502 CCNTINUE

KB2=1  
 IF(IP1,LT,K(JOJ))KB2=0  
 IF(IP1,GT,L(JOJ))KB2=0

KB3=1  
 IF(IM1,LT,K(JOJ))KB3=0  
 IF(IM1,GT,L(JOJ))KB3=0

IF(JPI,GT,NY) GO TO 503

C  
 C  
 C

C  
 C  
 C  
 C  
 C  
 C  
 C  
 C

C

C

C

INTP0550  
INTP0555  
INTP0560  
INTP0565  
INTP0570  
INTP0575  
INTP0580  
INTP0585  
INTP0590  
INTP0595  
INTP0600  
INTP0605  
INTP0610  
INTP0615  
INTP0620  
INTP0625  
INTP0630  
INTP0635  
INTP0640  
INTP0645  
INTP0650  
INTP0655  
INTP0660  
INTP0665  
INTP0670  
INTP0675  
INTP0680  
INTP0685  
INTP0690  
INTP0695  
INTP0700  
INTP0705  
INTP0710  
INTP0715  
INTP0720  
INTP0725

```

KB4=1
IF(IOI.LT.K(JP1))KB4=0
IF(IOI.GT.L(JP1))KB4=0
GC TC 504
503 KB4=0
C
C
C
C
      DETERMINE COEFFICIENTS OF LOCAL PARABOLIC FIT
504 IF(KB2.EQ.1) GC TC 24
D=0.0
IF(KP3.EQ.1)GC TC 27
B=0.0
GC TC 26
24 IF(KB3.EQ.1)GC TC 25
D=0.0
P=W(L4)-W(L3)
GC TC 26
25 B=0.5*(W(L4)-W(L2))
D=(W(L4)+W(L2)-2.0*W(L3))/2.0
GC TC 26
27 B=W(L3)-W(L2)
C
26 IF(KB1.EQ.1)GC TC 28
E=0.0
IF(KP4.EQ.1)GC TC 31
C=0.0
GC TC 30
28 IF(KB4.EQ.1)GC TC 29
E=0.0
C=W(L3)-W(L1)
GC TC 30
29 C=0.5*(W(L5)-W(L1))
E=(W(L1)+W(L5)-2.0*W(L3))/2.0
GC TC 30
31 C=W(L5)-W(L3)

```



INTP0910  
INTP0915  
INTP0920  
INTP0925  
INTP0930  
INTP0935  
INTP0940  
INTP0945  
INTP0950  
INTP0955  
INTP0960  
INTP0965  
INTP0970  
INTP0975  
INTP0980  
INTP0985  
INTP0990  
INTP0995  
INTP1000  
INTP1005  
INTP1010  
INTP1015  
INTP1020  
INTP1025  
INTP1030  
INTP1035  
INTP1040  
INTP1045  
INTP1050  
INTP1055  
INTP1060  
INTP1065  
INTP1070  
INTP1075  
INTP1080  
INTP1085

WWX(2)=WWX(4)  
GO TO 17  
183 WW(3)=0.5\*(WW(1)+WW(4))  
WWX(3)=0.5\*(WWX(1)+WWX(4))  
GO TO 17  
184 WW(1)=WW(2)  
WWX(1)=WWX(2)  
WW(3)=WW(4)  
WWX(3)=WWX(4)  
GO TO 17  
185 WW(3)=0.5\*(WW(1)+WW(4))  
WWX(3)=0.5\*(WWX(1)+WWX(4))  
WW(2)=WW(3)  
WWX(2)=WWX(3)  
GO TO 17  
186 WW(1)=WW(4)  
WW(2)=WW(4)  
WW(3)=WW(4)  
WWX(1)=WWX(4)  
WWX(2)=WWX(4)  
WWX(3)=WWX(4)  
GO TO 17  
187 WW(4)=0.5\*(WW(2)+WW(3))  
WWX(4)=0.5\*(WWX(2)+WWX(3))  
GO TO 17  
188 WW(1)=0.5\*(WW(2)+WW(3))  
WWX(1)=0.5\*(WWX(2)+WWX(3))  
WW(4)=WW(1)  
WWX(4)=WWX(1)  
GO TO 17  
189 WW(2)=WW(1)  
WWX(2)=WWX(1)  
WW(4)=WW(3)  
WWX(4)=WWX(3)  
GO TO 17  
190 WW(1)=WW(3)



```

WW(2)=WW(3)
WW(4)=WW(3)
WWX(1)=WWX(3)
WWX(2)=WWX(3)
WWX(4)=WWX(3)
GO TO 17
191 WW(3)=WW(1)
WWX(3)=WWX(1)
WW(4)=WW(2)
WWX(4)=WWX(2)
GO TO 17
192 WW(1)=WW(2)
WW(3)=WW(2)
WW(4)=WW(2)
WWX(1)=WWX(2)
WWX(3)=WWX(2)
WWX(4)=WWX(2)
GO TO 17
193 WW(2)=WW(1)
WW(3)=WW(1)
WW(4)=WW(1)
WWX(2)=WWX(1)
WWX(3)=WWX(1)
WWX(4)=WWX(1)
GO TO 17

```

```

C 101 WRITE(6,111)
WRITE(6,112) XO,YO
111 FORMAT(49HNO NEIGHBORING POINTS-NO INTERPOLATION ATTEMPTED)
112 FORMAT(1H0,32HCOORDINATES OF POINT IN QUESTION//25X,3HXO ,E17 8,3HINTP1235
1YO=,E27.8)
GO TO 13

```

```

C
C
C COEFFICIENTS OF WEIGHTING FUNCTION WANS
C

```

```

INTP1090
INTP1095
INTP1100
INTP1105
INTP1110
INTP1115
INTP1120
INTP1125
INTP1130
INTP1135
INTP1140
INTP1145
INTP1150
INTP1155
INTP1160
INTP1165
INTP1170
INTP1175
INTP1180
INTP1185
INTP1190
INTP1195
INTP1200
INTP1205
INTP1210
INTP1215
INTP1220
INTP1225
INTP1230
INTP1235
INTP1240
INTP1245
INTP1250
INTP1255
INTP1260
INTP1265

```

INTP1270  
INTP1275  
INTP1280  
INTP1285  
INTP1290  
INTP1295  
INTP1300  
INTP1305  
INTP1310  
INTP1315  
INTP1320  
INTP1325  
INTP1330  
INTP1335

```

17 AA=WW(1)
   AAX=WWX(1)
   BB=WW(2)-WW(1)
   BBX=WWX(2)-WWX(1)
   CC=WW(3)-WW(1)
   CCX=WWX(3)-WWX(1)
   DD=(WW(4)-WW(3))-(WW(2)-WW(1))
   DDX=(WWX(4)-WWX(3))-(WWX(2)-WWX(1))

C
18 WANS=AA+BB *XOX+CC*YDY+DD*XOX*YDY
   WXANS=AAX+BBX*XOX+CCX*YDY+DDX*XOX*YDY+((BB+DD*YDY)/DX)
13 CONTINUE
   RETURN
   END

```

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